

GENERIC PROPERTIES OF 3-DIMENSIONAL REEB FLOWS: BIRKHOFF SECTIONS AND ENTROPY

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ABSTRACT. In this paper we use broken book decompositions to study Reeb flows on closed 3-manifolds. We show that if the Liouville measure of a nondegenerate contact form can be approximated by periodic orbits, then there is a Birkhoff section for the associated Reeb flow. In view of Irie's equidistribution theorem, this is shown to imply that the set of contact forms whose Reeb flows have a Birkhoff section contains an open and dense set in the C^∞ -topology. We also show that the set of contact forms whose Reeb flows have positive topological entropy is open and dense in the C^∞ -topology.

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1. INTRODUCTION AND RESULTS

The notion of a Birkhoff section (see Definition 2.4) is a classical tool to study flows on 3-manifolds, going back to Poincaré. When a flow admits such a section, its dynamics can be reduced to the dynamics of the first-return map on the section – a much simpler data. Birkhoff sections are particularly convenient for studying topological and dynamical properties of the flow. In contact topology, work of Giroux [24] implies that every contact structure on a closed 3-manifold admits a supporting open book decomposition. In particular, every contact structure is defined by a contact form whose Reeb vector field has a global surface of section, i.e. an embedded Birkhoff section. However for a given contact form – hence a given Reeb vector field – deciding if Birkhoff sections exist is difficult.

Not every flow on a closed 3-manifold admits Birkhoff sections. For example a flow which is not a suspension and has no periodic orbits does not have one. Notice however that such a flow cannot be Reeb in view of the result of Taubes [47] asserting that every Reeb vector field on a closed 3-manifold has a periodic orbit.

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On the other hand, many Reeb flows do admit Birkhoff sections. A large class of examples is provided by a striking result due to Hofer, Wysocki and Zehnder [26] which contains, as a special case, the fact that Hamiltonian flows on 3-dimensional strictly convex compact energy levels possess disk-like global surfaces of section. As a consequence, Hofer, Wysocki and Zehnder proved that on such an energy level there are two or infinitely many periodic orbits. Another class consists of geodesic flows on closed, orientable and connected Riemannian 2-manifolds such that the curvature has a definite sign. When the curvature is everywhere negative, this is implied by a result of Fried [21] asserting that transitive Anosov flows have Birkhoff sections. When the curvature is everywhere positive this is covered by a classical result of Birkhoff [7].

Consider a nonsingular vector field X on a compact 3-manifold M tangent to ∂M . If $\gamma : \mathbb{R}/T\mathbb{Z} \rightarrow M$ is a periodic orbit of period $T > 0$, parametrized by the flow, then $\frac{\gamma_* \text{Leb}}{T}$ is an invariant Borel probability measure, where Leb denotes the Lebesgue measure on $\mathbb{R}/T\mathbb{Z}$. This measure is independent of the choice of period. The set of invariant Borel probability measures is a convex subset of the topological dual of the space of continuous real-valued functions on M equipped with the supremum norm. A Borel probability measure is said to be approximated by periodic orbits if it is the weak* limit of a sequence of finite convex combinations of measures induced by periodic orbits.

When M is oriented, and X is the Reeb vector field of a positive contact form λ with helicity $\text{vol}(\lambda) = \int_M \lambda \wedge d\lambda > 0$, then we say that the Liouville measure can be approximated by periodic orbits if $(\lambda \wedge d\lambda)/\text{vol}(\lambda)$ can be approximated by periodic orbits as above. A periodic orbit is nondegenerate if its transverse linearized Poincaré map has no root of unity as an eigenvalue. The contact form is nondegenerate if every periodic Reeb orbit is nondegenerate. Birkhoff sections and ∂ -strong Birkhoff sections are defined in Section 2; ∂ -strong Birkhoff sections are, in particular, also Birkhoff sections. Our main result reads as follows.

THEOREM 1.1. *If the Liouville measure of a nondegenerate contact form on a closed 3-manifold can be approximated by periodic Reeb orbits, then the Reeb flow admits a ∂ -strong Birkhoff section.*

A result due to Irie [34] asserts that the set of nondegenerate contact forms on a closed 3-manifold whose Liouville measures can be approximated by periodic orbits is residual, in particular dense, with respect to the C^∞ -topology. Together with Proposition 5.1 we obtain the following statement.

COROLLARY 1.2. *On any closed 3-manifold, the set of contact forms such that the Reeb flow admits a ∂ -strong Birkhoff section spanned by nondegenerate periodic orbits is open and dense in the C^∞ -topology.*

The main key ingredient of the proof of Theorem 1.1 is the existence result found in [10] for *broken book decompositions* carrying the Reeb vector field of a nondegenerate contact form on any closed 3-manifold. These are special kinds of foliations on the complement of a finite set of periodic orbits, the so-called binding orbits, and whose leaves are transverse to the Reeb vector field. One of the main applications of broken book decompositions obtained in [10] is the statement that every nondegenerate Reeb flow on a closed 3-manifold has either two or infinitely many periodic orbits, generalizing a previous result by Cristofaro-Gardiner, Hutchings and Pomerleano [11] for cases where the first Chern class of the contact structure is torsion. All these results rely on the machinery of Embedded Contact Homology as defined by Hutchings [33].

REMARK 1.3. A result of Sigmund [46] implies that the Liouville measure of an Anosov Reeb flow can be approximated by periodic orbits. This fact and Theorem 1.1 together prove that Anosov Reeb flows have Birkhoff sections. This is a special case of Fried’s results from [21].

While writing this article, we learned that similar results were independently obtained by Contreras and Mazzucchelli [13].

THEOREM 1.4 (Contreras and Mazzucchelli [13]). *The Reeb flow of a nondegenerate contact form on a closed 3-manifold admits a Birkhoff section provided that the stable and unstable manifolds of all hyperbolic periodic Reeb orbits intersect transversely.*

A contact form whose Reeb flow satisfies the assumptions of the above result will be called here strongly nondegenerate, and were said to satisfy the Kupka-Smale condition in [13]. Their argument takes the path suggested in [10], see Remark 1.6, which is different from the one followed in this article. Both ways still start from a common crucial input: the existence result for broken book decompositions from [10]. In Appendix C we sketch an argument using the strategy elaborated in [10] and prior to the present work as well as to [13], but which leads to a generic existence result only valid in C^1 -topology, see Theorem C.1. It relies on the machinery developed by Arnaud, Bonatti and Crovisier [4, 8] and involves the proof of the *lift Axiom* (Lemma C.2) for Reeb flows. In turn, the lift Axiom and the derived connecting lemma (Theorem C.3) also imply Theorem C.4: the set of transitive Reeb vector fields is C^1 -generic.

Broken book decompositions generalize the finite-energy foliations obtained for nondegenerate Reeb flows on the tight 3-sphere by Hofer, Wysocki and Zehnder in their pioneering work [27]. The analysis from [27] led to major breakthroughs in the study of Reeb dynamics, as well as in the development of pseudo-holomorphic curve theory in symplectic cobordisms. In particular, Hofer, Wysocki and Zehnder were able to prove that a strongly nondegenerate Reeb flow on the tight 3-sphere has two or infinitely many periodic orbits, leading them to first conjecture that the $2/\infty$ dichotomy must hold for all Reeb flows on the tight 3-sphere.

The literature is full of results with sufficient existence conditions for Birkhoff sections, going back a century. Let us comment on the problem of finding global sections spanned by *given* periodic orbits, with no genericity assumptions. The simplest example is Birkhoff’s theorem [7]. Closed geodesics on oriented surfaces determine immersed annuli on the unit tangent bundle, made of unit vectors along the closed geodesic pointing “to one of the sides”. This is called a Birkhoff annulus. It is proved in [7] that Birkhoff annuli over embedded closed geodesics on positively curved spheres are global surfaces of section for the geodesic flow. This result is interesting because it gives strong dynamical conclusions out of geometric assumptions that can be checked in concrete examples. The analogue of Birkhoff’s result in the context of convex energy levels is found in [28], where the periodic orbits that span disk-like global surfaces of section are characterized. The convexity assumption was first dropped in [30], eventually leading to the generalization of Birkhoff’s result to Reeb flows via pseudo-holomorphic curves in [32], where it is shown that certain linking assumptions on invariant measures, well-known to be sufficient for global sections of general flows ([29, Theorem 1.3]), can sometimes be reduced to linking assumptions on periodic orbits.

It is also interesting to try to measure the minimal complexity that a Birkhoff section of a Reeb flow can have. For instance, Fried showed that geodesic flows on hyperbolic surface always admit torus-like Birkhoff sections [21] and this was recently extended to hyperbolic 2-dimensional orbifolds [14, 15]. On the other

hand, van Koert constructed Reeb flows on S^3 without disk-like global surfaces of section [48].

Our second result concerns the topological entropy of Reeb flows in dimension three. It is already known that a nondegenerate Reeb vector field with zero topological entropy admits a Birkhoff section [10]. Perturbing the first-return map on the section, thanks to the work of Koropecki, Le Calvez and Nassiri [38], and of Le Calvez and Sambarino [39] that we adapt to diffeomorphisms of surfaces with boundary, we obtain the following statement.

THEOREM 1.5. *Let M be a closed 3-manifold and $\xi \subset TM$ be a co-orientable contact structure. The set of contact forms on M defining ξ such that the Reeb flow has positive topological entropy is open and dense in the set of all contact forms on M defining ξ , with respect to the C^∞ -topology.*

This result can be seen as a generalization of the fact, shown by Knieper and Weiss [37], that the geodesic flow on a closed surface has positive topological entropy C^∞ -generically on the metric. Our proof follows essentially the same path: they used the pseudo-holomorphic curves of Hofer, Wysocki and Zehnder [27] to produce Birkhoff sections, dealing with the extra constraint of perturbing the metric instead of the flow among its Reeb neighbors.

There are many results available in the literature providing conditions that force a Reeb flow to have positive topological entropy. In [10] it was proved that if the 3-manifold is not graphed – for example if it is hyperbolic – then every nondegenerate Reeb vector field has positive topological entropy. Symplectic topological methods for studying positive topological entropy of Hamiltonian systems have now been vastly used, going back to Polterovich [42], Frauenfelder and Schlenk [19] and others. In [1] a condition on the fractional Dehn twist coefficients of a supporting open book decomposition is given to guarantee that every Reeb vector field, possibly degenerate, for the given supported contact structure has positive topological entropy. Alves and Pirnapasov [3] proved that any contact 3-manifold admits transverse links that force positive topological entropy when realized as periodic Reeb orbits. Alves and Meiwes obtain in [2] contact structures in higher dimensional spheres such that the Reeb flows of all defining contact forms have positive topological entropy.

We end this introduction with a rough outline of the proof of Theorem 1.1.

Given a flow in a closed manifold M and given a class y in $H^1(M)$, Schwartzman [44, Section 7] gave a necessary and sufficient condition for the flow to admit a Birkhoff section with empty boundary dual to y : for every invariant Borel probability measure μ the intersection of μ with y , see Section 3.3 for a definition, must be positive. This criterion can be adapted to the existence of a Birkhoff section with prescribed boundary, via Schwartzman-Fried-Sullivan theory [20, 44, 45]. It was already known to specialists for some time but we provide an explicit statement as Theorem A.1: given a link L made of periodic orbits and a class y in $H^1(M \setminus L)$, there exists a ∂ -strong Birkhoff section if every invariant measure supported in $M \setminus L$ intersects y positively and on every component of L the flow winds positively with respect to y . A refinement for global surfaces of section can be found in [29].

It is shown in [10] that when a contact form is nondegenerate its Reeb vector field is supported by a broken book decomposition (K, \mathcal{F}) in a special manner. Here K is a link formed by periodic orbits, called binding orbits, and \mathcal{F} is a special smooth foliation of $M \setminus K$, see Definition 2.7 for a precise description. The binding splits into special sublinks $K = K_r \cup K_b$, the components K_b are the so-called broken binding orbits. It follows from definitions that when $K_b = \emptyset$, the broken book is a

rational open book decomposition whose pages are Birkhoff sections for the Reeb flow. Hence we proceed assuming that $K_b \neq \emptyset$.

In this case the foliation \mathcal{F} contains a finite set of special leaves, called rigid leaves, see Section 3.2, which almost meet the conditions of Theorem A.1: their union $S_{R_{\mathcal{F}}}$ intersects all orbits of the flow. However it is not true that the flow winds positively with respect to $S_{R_{\mathcal{F}}}$, see Figure 1 right. In order to enforce this property, one would like to find a surface transverse to the flow, bounded by periodic orbits in $M \setminus K_b$, and that intersects all components of K_b positively. Then one could “add it” to the surface $S_{R_{\mathcal{F}}}$, thus obtaining the desired Birkhoff section, using a process that can be traced back to Fried’s work [21], see also [10].

We do not find directly such a transverse surface. However we remark that it is enough to find a surface bounded by periodic orbits and which intersects the broken binding orbits positively — the point being that this surface need not be positively transverse to the flow in its interior as is required in Fried’s process. Indeed, denoting by S_b this surface, we remark that the “sum” $nS_{R_{\mathcal{F}}} + S_b$ intersects positively all invariant measures when n is large enough, and that the flow winds positively with respect to it along the broken binding orbits.

At the technical level, it is easier to find the Poincaré dual to this surface S_b : we will find an additional link of periodic orbits $K' \subset M \setminus K_b$ and a cohomology class $y' \in H^1(M \setminus K'; \mathbb{R})$ which essentially plays the role of the Poincaré dual of S_b . Once this task is done, Theorem 1.1 is a direct consequence of Theorem 2.10, stated in Section 2 and proved in Section 3. It uses the fine structural dynamical information revealed by the broken book decomposition in order to check the assumptions of Theorem A.1.

We are then left with the task of finding K' and y' , which is the content of Proposition 2.11. For this we need to establish a small modification of the Action-Linking Lemma from [6]. This lemma relates the intersection number between a given surface and the Liouville measure (a topological data) to the contact area of the surface. Hence, if the Liouville measure can be approximated by periodic orbits then the contact area is re-obtained as the limit of intersection numbers of weighted links of periodic orbits with the surface. The link K' is found among those links.

In the special case where M is a homology 3-sphere, we can make the argument explicit and short: denoting by I_k a sequence of weighted links that approximate the Liouville measure, and by h_1, \dots, h_n the broken binding orbits, the Action-Linking Lemma implies that the linking $\text{Lk}(h_i, \lambda \wedge d\lambda)$ equals the action $T(h_i)$, which is positive. Since $\text{Lk}(h_i, \lambda \wedge d\lambda)$ is the limit of the sequence $\text{Lk}(h_i, I_k)$ when $k \rightarrow \infty$, we obtain that for k large enough, all linking numbers $\text{Lk}(h_i, I_k)$, $i = 1, \dots, n$, are positive. This implies that I_k links positively with all broken binding orbits. Therefore we can take $K' = I_k$ and $y' = \text{Lk}(\cdot, I_k)$.

This argument is special to homology 3-spheres, in general it needs to be replaced by a refined one found in Section 4. Once these tools are in place, Theorem 1.1 has the following simple proof.

Proof of Theorem 1.1. By [10, Theorem 1.1] we know that the Reeb vector field of a nondegenerate contact form is strongly carried by a broken book decomposition, see Definition 2.7 and Remark 2.8. An application of Proposition 2.11 with $L = K_b$ tells us that the hypotheses of Theorem 2.10 are satisfied. \square

REMARK 1.6. Contreras and Mazzucchelli studied in [13] the closure of the invariant manifolds of the broken binding orbits, and found many homoclinic connections. Once enough such homoclinic connections are found, the strategy described in [10] and also implemented in the proof of Theorem C.1 applies. It is possible to show

that there are additional periodic orbits spanned by transverse surfaces that intersect the broken binding orbits. In [10, Theorem 4.13] it is then explained that these additional surfaces can be “added” to form a new broken book decomposition carried by the contact form, but with strictly less broken binding orbits. After a finite number of steps, one gets a rational open book decomposition whose pages are Birkhoff sections. A key contribution from [13], which seems surprising and of independent interest, is that under the strong nondegeneracy assumption the closure of stable and unstable manifolds of the broken binding orbits coincide. Using this fact, the desired additional transverse surfaces can be constructed and the strategy described above can be applied. At this point of the argument, these additional surfaces provide precisely the data K', y' and the reduction process can be replaced by our Theorem 2.10.

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2. BROKEN BOOK DECOMPOSITIONS

2.1. Rotation numbers along periodic orbits. Let M be any orientable smooth 3-manifold, and let X be a smooth vector field on M . Consider a periodic orbit $\gamma \subset M \setminus \partial M$ of X , with primitive period $T > 0$, as a knot oriented by X . This orientation and the ambient orientation together co-orient γ . Denote by ϕ^t the local flow near γ . Consider a small compact neighborhood N of γ , and an orientation preserving diffeomorphism

$$(1) \quad \Psi : N \rightarrow \mathbb{R}/T\mathbb{Z} \times \mathbb{D}$$

satisfying $\Psi(\phi^t(p_0)) = (t, 0)$ for some $p_0 \in \gamma$. Here the closed unit disk $\mathbb{D} \subset \mathbb{C}$ and the circle $\mathbb{R}/T\mathbb{Z}$ are oriented by the canonical orientations of \mathbb{C} and \mathbb{R} respectively, and $\mathbb{R}/T\mathbb{Z} \times \mathbb{D}$ gets the product orientation. On $N \setminus \gamma$ we have coordinates

$$(2) \quad (t, r, \theta) \in \mathbb{R}/T\mathbb{Z} \times (0, 1] \times \mathbb{R}/2\pi\mathbb{Z}, \quad \Psi^{-1}(t, re^{i\theta}) \simeq (t, r, \theta)$$

loosely referred to as tubular polar coordinates around γ . Consider the vector bundle $E_\gamma = TM|_\gamma/T\gamma \rightarrow \gamma$ oriented by the co-orientation of γ , and denote by E_γ^* the complement of its zero section. The total space of the circle bundle $E_\gamma^*/\mathbb{R}_+ \rightarrow \gamma$ is a torus, which can be equipped with global coordinates $(t, \theta) \in \mathbb{R}/T\mathbb{Z} \times \mathbb{R}/2\pi\mathbb{Z}$ induced by Ψ . The linearized flow $D\phi^t$ along γ descends to a flow on E_γ^*/\mathbb{R}_+ represented as the flow of a vector field of the form

$$(3) \quad \partial_t + b(t, \theta)\partial_\theta$$

on this torus. The smooth function $b(t, \theta)$ is $(T\mathbb{Z} \times 2\pi\mathbb{Z})$ -periodic.

DEFINITION 2.1. If $y \in H^1(N \setminus \gamma; \mathbb{R})$ is cohomologous to $p dt + q d\theta$, with constants $p, q \in \mathbb{R}$, then define the rotation number of γ relative to y as

$$(4) \quad \rho^y(\gamma) = \frac{T}{2\pi} \left(p + q \lim_{t \rightarrow +\infty} \frac{\theta(t)}{t} \right)$$

where $\theta : \mathbb{R} \rightarrow \mathbb{R}$ is any solution of $\dot{\theta}(t) = b(t, \theta(t))$.

REMARK 2.2. The rotation number $\rho^y(\gamma)$ does not depend on the choice of Ψ , or on the choice of the solution $\theta(t)$; see [29, Section 2] for details.

REMARK 2.3. If $\gamma \subset M \setminus \partial M$ as above is a component of some link $L \subset M$ and $y \in H^1(M \setminus L; \mathbb{R})$, then we may identify y with the element of $H^1(N \setminus \gamma; \mathbb{R})$ obtained by pulling y back via the inclusion map $N \setminus \gamma \hookrightarrow M \setminus L$. We still write $\rho^y(\gamma)$ for the corresponding rotation number, with no fear of ambiguity.

2.2. Broken books and Birkhoff sections. Consider a closed, connected, orientable and smooth 3-manifold M with a smooth nonsingular vector field X . The flow of X is denoted by ϕ^t .

DEFINITION 2.4. A section for the flow of X is an immersion $\iota : S \rightarrow M$ defined on a compact surface S such that:

- (i) If $\partial S \neq \emptyset$ then $\iota(\partial S)$ is a link consisting of periodic orbits of X .
- (ii) $\iota^{-1}(\iota(\partial S)) = \partial S$ and ι defines an embedding $S \setminus \partial S \hookrightarrow M \setminus \iota(\partial S)$ transverse to X .

A Birkhoff section, or rational global surface of section, for the flow of X is a section such that:

- (iii) For every $p \in M$ there exist $t_- < 0 < t_+$ such that $\phi^{t_\pm}(p) \in \iota(S)$.

If ι is simultaneously an embedding and a Birkhoff section, then $\iota(S)$ is called a global surface of section for the flow of X .

If γ is a periodic orbit of X then $E_\gamma = TM|_\gamma/T\gamma$ is a vector bundle over γ . Hence $\mathbb{P}^+\gamma = (E_\gamma \setminus 0)/\mathbb{R}_+ \rightarrow \gamma$ is a circle bundle. Note that the linearized flow $D\phi^t|_\gamma$ determines a smooth flow on $\mathbb{P}^+\gamma$. We call this flow the linearized flow on $\mathbb{P}^+\gamma$, with no fear of ambiguity. If $\iota : S \rightarrow M$ is a section for the flow of X and c is a connected component of ∂S such that $\iota(c) = \gamma$, then ι defines a smooth map $\nu_c^c : c \rightarrow \mathbb{P}^+\gamma$ as follows: choose any smooth vector field n of S along c pointing outwards, so that $d\iota(n)$ is a map $c \rightarrow TM|_\gamma$, and define ν_c^c to be the map obtained by composing $d\iota(n)$ with the quotient map $TM|_\gamma \rightarrow \mathbb{P}^+\gamma$. The definition of ν_c^c does not depend on the choice of n . The trace of ι along c is defined as the image of the map ν_c^c , in particular it is a subset of $\mathbb{P}^+\gamma$.

DEFINITION 2.5. Let $\iota : S \rightarrow M$ be a section for the flow of X .

- We call ι a ∂ -strong section if its trace along every connected component $c \subset \partial S$ is an embedding transverse to the linearized flow on $\mathbb{P}^+\gamma$; here γ is the periodic orbit in $\iota(\partial S)$ that contains $\iota(c)$.
- If ι is a Birkhoff section then we call ι a ∂ -strong Birkhoff section if it is a ∂ -strong section, and if for every connected component $c \subset \partial S$ its trace along c defines a global section for the linearized flow on $\mathbb{P}^+\gamma$; here γ is the periodic orbit that contains $\iota(c)$.

REMARK 2.6. When the ∂ -strong Birkhoff section is a global surface of section, then we get a ∂ -strong global surface of section as in [18, Definition 1.6].

DEFINITION 2.7 ([10]). Assume now that M is oriented. The vector field X is said to be strongly carried by a broken book decomposition (K, \mathcal{F}) , where the binding $K \subset M$ is a link consisting of periodic orbits and \mathcal{F} is a smooth foliation of $M \setminus K$, if the following conditions are satisfied.

- (I) The binding K , which is oriented by the flow, splits into two sublinks as $K = K_r \sqcup K_b$, where K_r is called the radial part of the binding and K_b the broken part of the binding. The leaves of \mathcal{F} are transverse to and co-oriented by X , and oriented by the orientation of M and this co-orientation. In particular, trajectories intersect the leaves positively.
- (II) For every leaf ℓ of \mathcal{F} there exists a compact connected oriented surface S with non-empty boundary, and a section $\iota : S \rightarrow M$ for the flow such that $\iota|_{S \setminus \partial S}$ defines an orientation preserving diffeomorphism $S \setminus \partial S \rightarrow \ell$, and $\iota|_{\partial S}$ defines a (not necessarily surjective) submersion $\partial S \rightarrow K$. Let c be a connected component of ∂S with $\iota(c) = \gamma \subset K$, in which case we say that γ is in the boundary of ℓ . Let $k \in \mathbb{Z} \setminus \{0\}$ be the degree of $\iota|_c : c \rightarrow \gamma$, and let $\ell^* \in H^1(M \setminus K; \mathbb{R})$ be the class dual to ℓ .
 - (a) If $\gamma \in K_r$ then $\rho^{\ell^*}(\gamma) > 0$.
 - (b) If $\gamma \in K_b$ then $|k| \in \{1, 2\}$, $\rho^{\ell^*}(\gamma) = 0$, γ is hyperbolic and its transverse linearized Poincaré map has real eigenvalues $\alpha < \beta$. If $|k| = 1$ then γ is positive hyperbolic in the sense that $0 < \alpha < 1 < \beta$. If $|k| = 2$ then γ is negative hyperbolic in the sense that $\alpha < -1 < \beta < 0$.
- (III) If $\gamma \subset K_r$ then the intersection of the leaves of \mathcal{F} with a small disk D transverse to γ defines a radial foliation of $D \setminus \gamma$ centered at $D \cap \gamma$.
- (IV) If $\gamma \subset K_b$ then the intersections of the leaves of \mathcal{F} with a small disk D transverse to γ divide $D \setminus \gamma$ into eight sectors centered at $D \cap \gamma$. Four of these sectors do not intersect $W^s(\gamma) \cup W^u(\gamma)$, are radially foliated and might have empty interior. These are intercalated by four open sectors containing $(W^s(\gamma) \cup W^u(\gamma)) \cap (D \setminus \gamma)$, which are foliated by hyperbola.

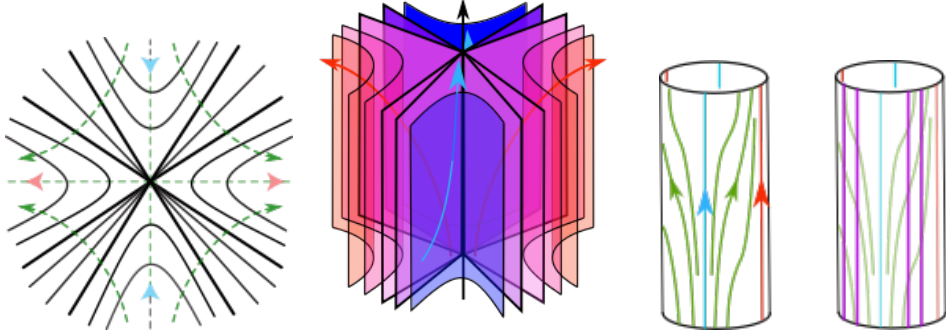


FIGURE 1. A broken book decomposition strongly carrying a vector field X in a neighbourhood of a broken binding component γ . On the left, the transversal view: the intersection of the broken book with a disc transverse to X . The special leaves are bolded. Some orbits of X are represented with green dotted lines. In particular the four local stable/unstable manifolds lie in four different hyperbolic sectors. In the center the 3d-view. On the right, the dynamics on the blown-up broken binding orbit $\mathbb{P}^+\gamma$ and the special leaves (purple): they are transverse to the projectivized flow and intersect all its orbits, except those corresponding to the stable and unstable manifolds. This manifests the fact that the rotation number of the flow along a broken binding component with respect to the union of the special leaves $\rho^{\ell^*}(\gamma)$ is zero.

REMARK 2.8. In [10] it is defined that a contact form is carried by the broken book decomposition (K, \mathcal{F}) if all the properties in Definition 2.7 hold for the Reeb

vector field X , except for (IIa). Without further assumptions, one gets a non-strict inequality $\rho^{\ell^*}(\gamma) \geq 0$ for all $\gamma \subset K_r$ and for all leaves ℓ . But if $\rho^{\ell^*}(\gamma) = 0$ for some $\gamma \subset K_r$ and some leaf ℓ that has γ on the boundary, then some iterate of γ is degenerate in the sense that 1 is an eigenvalue of the corresponding transverse Poincaré map. Hence, if λ is a nondegenerate contact form carried by (K, \mathcal{F}) as in [10], then its Reeb vector field is strongly carried by (K, \mathcal{F}) as in Definition 2.7.

REMARK 2.9. The finite-energy foliations obtained by Hofer, Wysocki and Zehnder in [27] for a nondegenerate Reeb flow on the standard contact 3-sphere are \mathbb{R} -invariant foliations of $\mathbb{R} \times S^3$ whose leaves are pseudo-holomorphic curves with finite Hofer energy, all of which have genus zero and precisely one positive puncture. They induce a broken book decomposition (K, \mathcal{F}) as above with several special properties. The \mathbb{R} -invariant cylinders over the orbits of K are precisely the \mathbb{R} -invariant leaves of the finite-energy foliation, the other leaves project to the leaves of \mathcal{F} . Moreover, all orbits in K have self-linking number equal to -1 , all orbits in K_b are positive hyperbolic, and each of the four radially foliated sectors in (IV) consists of a single ray. There are several other special additional properties.

THEOREM 2.10. *Suppose that X is strongly carried by a broken book decomposition $(K = K_b \sqcup K_r, \mathcal{F})$, and that there exists a link $K' \subset M \setminus K_b$ consisting of periodic orbits and a class $y' \in H^1(M \setminus K'; \mathbb{R})$ satisfying $\langle y', \gamma \rangle > 0$ for every $\gamma \subset K_b$. Then there is a ∂ -strong Birkhoff section for the flow of X with boundary in $K \cup K'$.*

Let λ be a positive contact form on M . Consider the space $C^0(M)$ of continuous real-valued functions on M . It becomes a Banach space with the supremum norm. As described in the introduction, we shall say that the Liouville measure can be approximated by periodic orbits if there exists a sequence of weighted links of periodic Reeb orbits

$$(\{\gamma_j^n\}, \{p_j^n\}) \quad j = 1, \dots, N(n)$$

where each γ_j^n is a periodic Reeb orbit with primitive period $T(\gamma_j^n) > 0$ and the weights $p_j^n \in (0, 1]$ satisfy $\sum_j p_j^n = 1$, such that

$$\lim_{n \rightarrow \infty} \sum_j p_j^n \frac{(\gamma_j^n)_* \text{Leb}}{T(\gamma_j^n)} = \frac{\lambda \wedge d\lambda}{\text{vol}(\lambda)}.$$

Here the orbits are seen as maps $\gamma_j^n : \mathbb{R}/T(\gamma_j^n)\mathbb{Z} \rightarrow M$ parametrized by the flow, Leb denotes Lebesgue measure on $\mathbb{R}/T(\gamma_j^n)\mathbb{Z}$, $\text{vol}(\lambda) = \int_M \lambda \wedge d\lambda$, and the limit is taken in the weak* topology of the topological dual $C^0(M)'$ of $C^0(M)$.

PROPOSITION 2.11. *Suppose that X is the Reeb vector field of a positive contact form on M such that $\lambda \wedge d\lambda$ can be approximated by periodic Reeb orbits. Let $L \subset M$ be a link consisting of periodic Reeb orbits. Then there exists a link $K' \subset M \setminus L$ and $y' \in H^1(M \setminus K'; \mathbb{R})$ such that K' consists of periodic Reeb orbits, and $\langle y', \gamma \rangle > 0$ for every $\gamma \subset L$.*

3. FROM LINKING WITH THE BROKEN BINDING TO BIRKHOFF SECTIONS

The goal of this section is to prove Theorem 2.10. Throughout this section we consider a broken book decomposition (K, \mathcal{F}) strongly carrying a smooth vector field X on a closed, connected and oriented 3-manifold M as in Definition 2.7. The flow of X is denoted by ϕ^t .

Note that $K_b = \emptyset$ if, and only if, the broken book decomposition is a rational open book decomposition as defined in [5], whose pages are Birkhoff sections. Hence we proceed assuming that $K_b \neq \emptyset$.

3.1. Blowing periodic orbits up. The discussion here makes no use of the broken book decomposition. We consider any link $L \subset M$ made of periodic orbits of ϕ^t , oriented by the flow. For each orbit $\gamma \subset L$ denote by $T_\gamma > 0$ its primitive period, and let N_γ be a small tubular neighborhood of γ with an orientation preserving diffeomorphism $\Psi_\gamma : N_\gamma \rightarrow \mathbb{R}/T_\gamma\mathbb{Z} \times \mathbb{D}$ as in (1). On $N_\gamma \setminus \gamma$ we have tubular polar coordinates $(t, r, \theta) \in \mathbb{R}/T_\gamma\mathbb{Z} \times (0, 1] \times \mathbb{R}/2\pi\mathbb{Z}$, $\Psi_\gamma^{-1}(t, re^{i\theta}) \simeq (t, r, \theta)$ as in (2). One can blow the link L up to construct a smooth 3-manifold M_L defined as

$$(5) \quad M_L := \left\{ M \setminus L \sqcup \bigsqcup_{\gamma \subset L} \mathbb{R}/T_\gamma\mathbb{Z} \times (-\infty, 1] \times \mathbb{R}/2\pi\mathbb{Z} \right\} / \sim$$

where (t, r, θ) in $\mathbb{R}/T_\gamma\mathbb{Z} \times (0, 1] \times \mathbb{R}/2\pi\mathbb{Z}$ gets identified with $\Psi_\gamma^{-1}(t, re^{i\theta})$. On M_L there are $\#\pi_0(L)$ smoothly embedded tori

$$(6) \quad \Sigma_\gamma = \mathbb{R}/T_\gamma\mathbb{Z} \times \{0\} \times \mathbb{R}/2\pi\mathbb{Z}$$

with coordinates $(t, \theta) \in \mathbb{R}/T_\gamma\mathbb{Z} \times \mathbb{R}/2\pi\mathbb{Z}$, and M_L contains a smooth compact domain

$$(7) \quad D_L = \left\{ M \setminus L \sqcup \bigsqcup_{\gamma \subset L} \mathbb{R}/T_\gamma\mathbb{Z} \times [0, 1] \times \mathbb{R}/2\pi\mathbb{Z} \right\} / \sim$$

satisfying

$$\partial D_L = \bigsqcup_{\gamma \subset L} \Sigma_\gamma.$$

Note that $D_L \setminus \partial D_L = M \setminus L$. Note also that $(t, r, \theta) \mapsto \Psi_\gamma^{-1}(t, re^{i\theta})$ defines a smooth diffeomorphism from $\mathbb{R}/T_\gamma\mathbb{Z} \times (0, 1] \times \mathbb{R}/2\pi\mathbb{Z}$ to $N_\gamma \setminus \gamma$.

It follows from arguments originally due to Fried [20] that X can be smoothly extended from $M \setminus L$ to a vector field X_L on M_L , whose flow we denote by ϕ_L^t . The restriction of X_L to D_L is unique, and X_L is tangent to ∂D_L . The details we need on Fried's construction can be found below, for more details see [29, Section 3]. In particular, it is important to know that there is a precise relation between the dynamics of X_L on Σ_γ and the linearized dynamics of X along γ . To see this, denote by $Z = (\Psi_\gamma)_* X$ the representation of X in $\mathbb{R}/T_\gamma\mathbb{Z} \times \mathbb{D}$. By the fundamental theorem of calculus we get that

$$Z(t, 0) = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \Rightarrow \quad Z(t, re^{i\theta}) = \begin{pmatrix} 1 \\ 0 \end{pmatrix} + A(t, re^{i\theta})re^{i\theta}$$

where

$$A(t, re^{i\theta}) = \int_0^1 D_2 Z(t, \tau re^{i\theta}) d\tau = \begin{pmatrix} A_1(t, re^{i\theta}) \\ A_2(t, re^{i\theta}) \end{pmatrix}$$

and D_2 stands for the partial derivative in the \mathbb{D} -factor. Note also that

$$D_2 Z(t, 0) = \begin{pmatrix} A_1(t, 0) \\ A_2(t, 0) \end{pmatrix}.$$

Now consider the smooth map

$$\Phi : \mathbb{R}/T_\gamma\mathbb{Z} \times [0, 1] \times \mathbb{R}/2\pi\mathbb{Z} \rightarrow \mathbb{R}/T_\gamma\mathbb{Z} \times \mathbb{D} \quad (t, r, \theta) \mapsto (t, re^{i\theta}).$$

Note that Φ defines a diffeomorphism $\mathbb{R}/T_\gamma\mathbb{Z} \times (0, 1] \times \mathbb{R}/2\pi\mathbb{Z} \simeq \mathbb{R}/T_\gamma\mathbb{Z} \times (\mathbb{D} \setminus \{0\})$ that can be used to pull $Z|_{\mathbb{R}/T_\gamma\mathbb{Z} \times (\mathbb{D} \setminus \{0\})}$ back to a vector field W . Since

$$D\Phi^{-1}(t, re^{i\theta}) = \begin{pmatrix} 1 & & 0 \\ 0 & \begin{pmatrix} \cos \theta & \sin \theta \\ -r^{-1} \sin \theta & r^{-1} \cos \theta \end{pmatrix} & \end{pmatrix}$$

it follows that

$$\begin{aligned} W(t, r, \theta) &= D\Phi^{-1}(t, re^{i\theta})Z(t, re^{i\theta}) \\ &= (1 + A_1(t, re^{i\theta})re^{i\theta})\partial_t + \langle e^{i\theta}, A_2(t, re^{i\theta})re^{i\theta} \rangle \partial_r + \langle ie^{i\theta}, A_2(t, re^{i\theta})e^{i\theta} \rangle \partial_\theta \end{aligned}$$

extends smoothly to $\mathbb{R}/T_\gamma\mathbb{Z} \times [0, 1] \times \mathbb{R}/2\pi\mathbb{Z}$. Since at $r = 0$ the component in ∂_r vanishes, we conclude that the extension X_L of X from $M \setminus L$ to D_L is tangent to ∂D_L . The restriction of X_L to Σ_γ is then of the form

$$(8) \quad \partial_t + b(t, \theta)\partial_\theta \quad b(t, \theta) = \langle ie^{i\theta}, A_2(t, 0)e^{i\theta} \rangle.$$

With the above formula, the relation between the dynamics of X_L on ∂D_L and the linearized dynamics along the various orbits $\gamma \subset L$ becomes precise, since $b(t, \theta)$ is exactly the same as the function appearing in (3). A solution

$$\begin{pmatrix} a(t) \\ u(t) \end{pmatrix} = D\phi^t(t_0, 0) \begin{pmatrix} a_0 \\ u_0 \end{pmatrix}, \quad u_0 = |u_0|e^{i\theta_0} \neq 0$$

along γ must satisfy

$$(9) \quad u(t) = |u(t)|e^{i\theta(t)} \quad \dot{\theta}(t) = b(t + t_0, \theta(t)), \quad \theta(0) = \theta_0.$$

3.2. Rigid leaves. For a broken book decomposition (K, \mathcal{F}) , the set of rigid leaves $R_{\mathcal{F}}$ was considered in [10]. When $K_b \neq \emptyset$ it can be alternatively defined as a smallest collection of leaves R satisfying the following property:

- (*) If U is a connected open set such that $\bar{U} \cap (K \cup (\cup_{\ell \in R} \ell)) = \emptyset$ then $\mathcal{F}|_U$ is defined by a fibration $U \rightarrow \mathbb{R}$.

For convenience of the reader, we present here a proof of the following result.

PROPOSITION 3.1. *A broken book decomposition with $K_b \neq \emptyset$ that is strongly carrying a vector field, has at least one and at most finitely many rigid leaves.*

In order to prove this result, we will use a global version of Reeb stability for foliations.

THEOREM 3.2 (Corollary 4.2.2, Chapter V [25]). *Let X be a C^0 compact manifold with nonempty boundary endowed with a transversely oriented C^0 -foliation \mathcal{F} that is tangent to ∂X . If all the leaves of \mathcal{F} are compact then X fibers over an interval and the foliation is given by the fibers.*

To study rigid leaves and prove Proposition 3.1, we first select finitely many special leaves of \mathcal{F} such that following holds:

- For every broken binding orbit $\gamma \subset K_b$, the leaves separating the four radially foliated sectors from the sectors foliated by hyperbolas in the local model described in (IV) are in the selected leaves. These are the leaves marked in bold in Figure 1 left.
- Every binding orbit $\gamma \subset K$ is in the boundary of one of the selected leaves.

We will denote this finite set of leaves by \mathcal{S} and call them the special leaves.

Proof of Proposition 3.1. It is not difficult to see that it suffices to show that some nonempty finite collection of leaves satisfies property (*).

The foliation \mathcal{F} of $M \setminus K$ is transversely oriented by the vector field X . This is (I) Definition 2.7.

Since every $\gamma \subset K_b$ is hyperbolic, any neighborhood γ contains a compact neighborhood N_γ of γ with the following properties. Assume first that γ is positive hyperbolic. Then:

- (i) N_γ is homeomorphic to a solid torus.
- (ii) ∂N_γ is decomposed into eight closed annuli with mutually disjoint interiors. Each of these annuli is smoothly embedded in M .

- (iii) Four of the annuli in (ii) are transverse to X , the flow enters N_γ along two of them and exists N_γ along the other two. Each of these four annuli is contained in a leaf of \mathcal{F} not in \mathcal{S} .
- (iv) The four annuli in (iii) are intercalated by four annuli tangent to X .
- (v) In the local model described in (IV) Definition 2.7, the leaves of \mathcal{F} corresponding to rays intersect transversely the four annuli tangent to X , and define a foliation by circles of four smaller annuli (possibly degenerated to a circle).

If γ is negative hyperbolic then there is an analogous decomposition of ∂N_γ into four closed and smoothly embedded annuli with mutually disjoint interiors, two are transverse to X and contained in leaves of \mathcal{F} not in \mathcal{S} , and the other two are tangent to X . The annuli tangent to X contain smaller annuli smoothly foliated by circles (possibly only one circle) that are intersections with leaves of \mathcal{F} corresponding to rays in the local model described in (IV) Definition 2.7. In any case, we denote by \mathcal{S}' the (finite) set of leaves of \mathcal{F} that, for some $\gamma \subset K_b$, contain one of the annuli in ∂N_γ transverse to X .

Using (III) Definition 2.7, any $\gamma \subset K_r$ admits a small neighborhood N_γ diffeomorphic to a solid torus such that ∂N_γ is smoothly foliated by circles made of transverse intersections of the leaves of \mathcal{F} with ∂N_γ . At least one of such circles is contained in a leaf in \mathcal{S} .

Consider the sets

$$N_K = \bigcup_{\gamma \subset K} N_\gamma \quad \Omega = M \setminus \left(N_K \cup \bigcup_{\ell \in \mathcal{S} \cup \mathcal{S}'} \ell \right).$$

Note that Ω is open. Cutting $M \setminus \overset{\circ}{N}_K$ along leaves of $\mathcal{S} \cup \mathcal{S}'$ we see that each connected component ω of Ω is the interior of a compact C^0 3-manifold with boundary Y with the following properties. We can decompose $\partial Y = \partial_\tau Y \cup \partial_\partial Y$, where $\partial_\tau Y$ corresponds to pieces of leaves in $\mathcal{S} \cup \mathcal{S}'$, and $\partial_\partial Y$ are annuli in ∂N_K foliated by circles consisting of intersections with leaves of \mathcal{F} . We now want to get rid of $\partial_\partial Y$. If $\partial_\partial Y \neq \emptyset$ we can take Y' to be a copy of Y with the opposite orientation, and paste Y and Y' along $\partial_\partial Y \cup \partial_\partial Y'$ in such a way to obtain a compact C^0 3-manifold $Y \cup_\partial Y'$ with boundary, endowed with a C^0 foliation tangent to its boundary. This foliation is transversely orientable, and all its leaves are compact. By Theorem 3.2, we obtain that $Y \cup_\partial Y'$ fibers over a closed interval and the foliation is given by the fibers. Thus for every open connected set W compactly contained in ω the foliation $\mathcal{F}|_W$ is defined by a fibration $W \rightarrow \mathbb{R}$. We conclude by noting that N_K can be taken arbitrarily small, and $\mathcal{S} \cup \mathcal{S}'$ is finite and nonempty. \square

Hence, assuming that $K_b \neq \emptyset$, the set $R_{\mathcal{F}}$ is finite, and on each connected open set U compactly contained in $M \setminus R_{\mathcal{F}}$ the foliation $\mathcal{F}|_U$ is defined by a fibration over \mathbb{R} . Such a fibration gives a function that is strictly increasing or strictly decreasing along the segments of orbits of X contained in U . Without loss of generality, we assume that the function is strictly increasing.

LEMMA 3.3. *Fix any metric on M . For every $\delta > 0$ there exists $T_\delta > 0$ such that the following holds: if $p \in M \setminus K$ satisfies $\text{dist}(\phi^{[-T_\delta, T_\delta]}(p), K_b) \geq \delta$, then there exists $t_* \in [-T_\delta, T_\delta]$ such that $\phi^{t_*}(p)$ belongs to some leaf in $R_{\mathcal{F}}$.*

Proof. We give a proof by contradiction. Assume that there exists some $\delta > 0$ such that the conclusion of the lemma is false. Hence there is a sequence of points $p_n \in M \setminus K$ and a sequence $t_n \in \mathbb{R}$ tending to infinity as $n \rightarrow \infty$, such that $\text{dist}(\phi^{[-t_n, t_n]}(p_n), K_b) \geq \delta$ and $\phi^{[-t_n, t_n]}(p_n) \cap (\cup_{\ell \in R_{\mathcal{F}}} \ell) = \emptyset$. Modulo passing to a subsequence, p_n converges to some point p_∞ such that $\text{dist}(\phi^{\mathbb{R}}(p_\infty), K_b) \geq \delta$. Consider the compact set $E = K \cup (\cup_{\ell \in R_{\mathcal{F}}} \ell)$.

We claim that $\text{dist}(\phi^{\mathbb{R}}(p_\infty), E) > 0$. First suppose, by contradiction, that $\phi^{\mathbb{R}}(p_\infty)$ contains points arbitrarily close to K_r . It follows that the trajectories of p_n spend arbitrarily large amounts of time arbitrarily near K_r , and by (II.a) in Definition 2.7, $\phi^{\mathbb{R}}(p_n)$ intersects some rigid leaf for n large enough. This contradiction shows that $\text{dist}(\phi^{\mathbb{R}}(p_\infty), K) > 0$. Again by contradiction, we assume that $\phi^{\mathbb{R}}(p_\infty)$ contains points arbitrarily close to $\cup_{\ell \in R_{\mathcal{F}}} \ell$. Using the transversality of the flow with the leaves and the fact that $\phi^{\mathbb{R}}(p_\infty)$ is at a positive distance to K , we find an intersection between $\phi^{\mathbb{R}}(p_\infty)$ and some rigid leaf. As before, this implies that $\phi^{[-t_n, t_n]}(p_n)$ intersects some rigid leaf, an absurd. We are done with the proof of $\text{dist}(\phi^{\mathbb{R}}(p_\infty), E) > 0$.

It follows that the ω -limit set of p_∞ satisfies $\text{dist}(\omega(p_\infty), E) > 0$, hence $\omega(p_\infty)$ is a compact subset of some connected open set U compactly contained in $M \setminus E$. Choose a recurrent point $q \in \omega(p_\infty)$. By Proposition 3.1 there is a continuous function $f : U \rightarrow \mathbb{R}$ strictly increasing along the segments of orbits in U , in particular along the orbit of q . This leads to a contradiction: a continuous function cannot be strictly increasing along a recurrent orbit. \square

3.3. Invariant measures and intersection numbers. As in the previous section, we consider a link $L \subset M$ tangent to X , with the orientation induced by X . We use the notation and the constructions from Section 3.1.

Denote by $\mathcal{P}_\phi(M \setminus L)$ the set of ϕ^t -invariant Borel probability measures on $M \setminus L$, and by $\mathcal{P}_\phi(D_L)$ the set of ϕ_L^t -invariant Borel probability measures on D_L . If $\mu \in \mathcal{P}_\phi(M \setminus L)$ and $y \in H^1(M \setminus L; \mathbb{R})$ then we may choose a closed 1-form β on $M \setminus L$ representing y such that $\iota_X \beta$ is bounded, and define

$$(10) \quad \mu \cdot y = \int_{M \setminus L} \iota_X \beta \, d\mu$$

called the intersection number of μ and y . The existence of β as above can be seen with the help of polar tubular coordinates, and is implicitly contained in Section 3.1, see [29, Section 2.1] for more details. The independence on the choice of β can be seen by showing that there exists $f_{\mu, y} \in L^1(\mu)$ and a Borel set $E \subset M \setminus L$ such that:

- $\mu(E) = 1$ and all points in E are recurrent.
- If $p \in E$ and $V \subset M \setminus L$ is an open contractible neighborhood of p , then $T_n^{-1} \langle y, k(T_n, p) \rangle \rightarrow f_{\mu, y}(p)$ as $n \rightarrow \infty$, where $T_n \rightarrow +\infty$ is any sequence such that $\phi^{T_n}(p) \rightarrow p$, and $k(T_n, p)$ are choices of loops obtained by concatenating to $\phi^{[0, T_n]}(p)$ a path from $\phi^{T_n}(p)$ to p inside V .
- The identity

$$\mu \cdot y = \int_{M \setminus L} f_{\mu, y} \, d\mu$$

holds.

These facts can be proved with a direct application of the ergodic theorem, for more details see [29, Section 2.1].

Similarly, for any $\mu \in \mathcal{P}_\phi(D_L)$ and $y \in H^1(D_L; \mathbb{R})$ one defines

$$(11) \quad \mu \cdot y = \int_{D_L} \iota_{X_L} \beta \, d\mu$$

where β is a closed 1-form on D_L representing y . Independence of β is easier in this case since D_L is compact.

We may freely identify $H^1(D_L; \mathbb{R}) \simeq H^1(M \setminus L; \mathbb{R})$ via pull-back by the inclusion map $M \setminus L = D_L \setminus \partial D_L \hookrightarrow D_L$, as in the statement below.

LEMMA 3.4 ([29], Lemma 3.6). *If γ is a component of L , $\mu \in \mathcal{P}_\phi(D_L)$ satisfies $\text{supp}(\mu) \subset \Sigma_\gamma$, and y is a class in $H^1(D_L; \mathbb{R})$, then $\mu \cdot y = \frac{2\pi}{T_\gamma} \rho^y(\gamma)$, where $T_\gamma > 0$ is the primitive period of γ .*

3.4. First estimates of intersection numbers. The cohomology class

$$(12) \quad y_0 = \sum_{\ell \in R_{\mathcal{F}}} \ell^* \in H^1(M \setminus K; \mathbb{R}) \simeq H^1(D_L; \mathbb{R}),$$

dual to the rigid leaves, plays a key role in our arguments.

LEMMA 3.5. *If $\mu \in \mathcal{P}_\phi(M \setminus K)$ is ergodic then $\mu \cdot y_0 > 0$.*

Proof. Any $\gamma \subset K_b$ is a hyperbolic orbit, hence isolated as an invariant set. Moreover, by hyperbolicity and condition (IV), we can find an isolating compact neighborhood N_γ of γ inside any neighborhood of γ fixed *a priori*, and some $\eta > 0$ depending on N_γ , such that the following holds: If $p \in M \setminus (K \cup W^u(\gamma) \cup W^s(\gamma))$ and if $\omega \neq \emptyset$ is a connected component of $\{t \in \mathbb{R} \mid \phi^t(p) \in N_\gamma\}$, then $\omega = [a, b]$ is a compact interval with non-empty interior, $\phi^{[a-\eta, a) \cup (b, b+\eta]}(p) \cap N_\gamma = \emptyset$, and $\exists t_* \in \omega$ such that $\phi^{t_*}(p)$ belongs to some leaf in $R_{\mathcal{F}}$.

Denote $W^u(K_b) = \bigcup_{\gamma \subset K_b} W^u(\gamma)$ and $W^s(K_b) = \bigcup_{\gamma \subset K_b} W^s(\gamma)$. Taking unions of various small N_γ as above, we find an isolating compact neighborhood N_b for K_b inside any neighborhood of K_b fixed *a priori*, and some $\eta > 0$ depending on N_b , with the following property:

- (*) If $p \in M \setminus (K \cup W^u(K_b) \cup W^s(K_b))$ and ω is a non-empty connected component of $\{t \in \mathbb{R} \mid \phi^t(p) \in N_b\}$, then $\omega = [a, b]$ is a compact interval with non-empty interior, $\phi^{[a-\eta, a) \cup (b, b+\eta]}(p) \cap N_b = \emptyset$, and $\exists t_* \in \omega$ such that $\phi^{t_*}(p) \in \bigcup_{\ell \in R_{\mathcal{F}}} \ell$.

Fix $\varphi : M \rightarrow [0, 1]$ continuous such that φ is identically equal to 1 near K_b , and is supported on a neighborhood U of K_b satisfying $\mu(U \setminus K_b) < \frac{1}{2}$. We can proceed assuming, without loss of generality, that $N_b \subset \varphi^{-1}(1)$. Since μ is ergodic, the following hold simultaneously for μ -almost all points $p \in M \setminus K$:

$$(13) \quad \begin{aligned} & \limsup_{T \rightarrow +\infty} \frac{1}{T} \text{Leb}(\{t \in [0, T] \mid \phi^t(p) \in N_b\}) \\ & \leq \lim_{T \rightarrow +\infty} \frac{1}{T} \int_0^T \varphi(\phi^t(p)) dt = \int_{M \setminus K} \varphi d\mu < \frac{1}{2}, \end{aligned}$$

$$(14) \quad \mu \cdot y_0 = \lim_{T \rightarrow +\infty} \frac{1}{T} \#\{t \in [0, T] \mid \phi^t(p) \in \ell \text{ for some } \ell \in R_{\mathcal{F}}\}.$$

From now on we fix such p , and denote

$$J = \{t \in \mathbb{R} \mid \phi^t(p) \in N_b\}.$$

If we fix an auxiliary metric on M , and choose $\delta > 0$ such that $\text{dist}(\cdot, K_b) > \delta$ on $M \setminus N_b$, then by Lemma 3.3 we find $T_* > 0$ such that if $\phi^{[c, d]}(p) \subset M \setminus N_b$ and $d - c \geq T_*$ then $\phi^{[c, d]}(p)$ intersects some leaf in $R_{\mathcal{F}}$.

By (*), for every $T > 0$ we write $J \cap [0, T]$ is a finite union of $n(T)$ maximal compact intervals, which are all at least η apart from each other. We list these intervals as $h_1, \dots, h_{n(T)}$ in an increasing fashion: $\sup h_j \leq \inf h_{j+1}$. It follows that $\inf h_{j+1} - \sup h_j \geq \eta$. Note that only h_1 and $h_{n(T)}$ may not be connected components of J .

Consequently, we can write $[0, T] \setminus J$ as a finite union of $m(T)$ maximal intervals which are open relatively to $[0, T]$, with $|n(T) - m(T)| \leq 1$, all of which have length at least equal to η . Denote by $\omega_1^+, \dots, \omega_{m^+}^+$ those intervals with length $\geq T_*$.

Denote by $\omega_1^-, \dots, \omega_{m_-(T)}^-$ those intervals with length $< T_*$. We have $m(T) = m_+(T) + m_-(T)$ and

$$[0, T] \setminus J = \omega_1^+ \cup \dots \cup \omega_{m_+(T)}^+ \cup \omega_1^- \cup \dots \cup \omega_{m_-(T)}^-.$$

Denote $H(T) = \#\{t \in [0, T] \mid \phi^t(p) \in \cup_{\ell \in R_{\mathcal{F}}} \ell\}$, i.e. $H(T)$ is the number of hitting times between $\phi^{[0, +\infty)}(p)$ and the union of the union of leaves in $R_{\mathcal{F}}$ up to time T . By (14) we have $T^{-1}H(T) \rightarrow \mu \cdot y_0$ as $T \rightarrow +\infty$. Our constructions so far imply that there is at least one hitting time on each $\omega_1^+, \dots, \omega_{m_+(T)}^+$ and, by (*), at least one hitting time on each $h_2, \dots, h_{n(T)-1}$. In particular, $H(T) \geq n(T) - 2$. We split the remaining arguments in two cases.

Case 1. $\liminf_{T \rightarrow +\infty} T^{-1}n(T) > 0$

In this case we are done since

$$\mu \cdot y_0 = \lim_{T \rightarrow +\infty} \frac{H(T)}{T} \geq \liminf_{T \rightarrow +\infty} \frac{n(T)}{T} > 0.$$

Case 2. $\liminf_{T \rightarrow +\infty} T^{-1}n(T) = 0$

In this case we can find $T_j \rightarrow +\infty$ such that $T_j^{-1}n(T_j) \rightarrow 0$ as $j \rightarrow +\infty$. Note that

$$\begin{aligned} 1 &= \frac{\text{Leb}(J \cap [0, T_j])}{T_j} + \frac{\text{Leb}([0, T_j] \setminus J)}{T_j} \\ (15) \quad &= \frac{\text{Leb}(J \cap [0, T_j])}{T_j} + \frac{1}{T_j} \sum_{s=1}^{m_+(T_j)} \text{Leb}(\omega_s^+) + \frac{1}{T_j} \sum_{s=1}^{m_-(T_j)} \text{Leb}(\omega_s^-). \end{aligned}$$

Moreover, if $j \gg 1$ then $T_j^{-1}\text{Leb}(J \cap [0, T_j]) < \frac{1}{2}$ by (13), and

$$\frac{1}{T_j} \sum_{s=1}^{m_-(T_j)} \text{Leb}(\omega_s^-) \leq T_* \frac{m_-(T_j)}{T_j} \leq T_* \frac{m(T_j)}{T_j} \leq T_* \frac{n(T_j) + 1}{T_j} \rightarrow 0.$$

Plugging in (15) we get

$$(16) \quad j \gg 1 \quad \Rightarrow \quad \frac{1}{T_j} \sum_{s=1}^{m_+(T_j)} \text{Leb}(\omega_s^+) \geq \frac{1}{2}.$$

Consider λ_j the maximal number of intervals of length T_* that fit inside the union of the ω_s^+ . Note that

$$\left| T_* \lambda_j - \sum_{s=1}^{m_+(T_j)} \text{Leb}(\omega_s^+) \right| \leq T_* m_+(T_j).$$

Hence for j large enough we compute

$$\frac{1}{2} \leq \frac{1}{T_j} \sum_{s=1}^{m_+(T_j)} \text{Leb}(\omega_s^+) \leq \frac{T_* \lambda_j}{T_j} + \frac{T_* m_+(T_j)}{T_j}.$$

Since $\frac{m_+(T_j)}{T_j} \rightarrow 0$ we get $\frac{\lambda_j}{T_j} \geq \frac{1}{2T_*}$ for all j large enough. Hence

$$\frac{H(T_j)}{T_j} \geq \frac{\lambda_j}{T_j} \geq \frac{1}{2T_*}$$

for all j large enough, and $\mu \cdot y_0 \geq \frac{1}{2T_*} > 0$ follows. \square

3.5. Further estimates of intersection numbers. Let the broken book decomposition $(K = K_r \sqcup K_b, \mathcal{F})$, the link of periodic orbits $K' \subset M \setminus K_b$ and the cohomology class $y' \in H^1(M \setminus K'; \mathbb{R})$ be as in the statement of Theorem 2.10. As in Section 3.1, we may blow $K \cup K'$ up to obtain a new manifold $M_{K \cup K'}$ with a smooth compact domain $D_{K \cup K'}$ such that $D_{K \cup K'} \setminus \partial D_{K \cup K'} = M \setminus (K \cup K')$. There is a boundary torus Σ_γ (6) for each orbit $\gamma \subset K \cup K'$ and the vector field X extends smoothly from $M \setminus (K \cup K')$ to $D_{K \cup K'}$. As agreed before, we may freely identify $H^1(M \setminus (K \cup K'); \mathbb{R})$ with $H^1(D_{K \cup K'}; \mathbb{R})$.

We pull y' back to a cohomology class

$$(17) \quad y'' \in H^1(M \setminus (K \cup K'); \mathbb{R}) \simeq H^1(D_{K \cup K'}; \mathbb{R})$$

via the inclusion $M \setminus (K \cup K') \hookrightarrow M \setminus K'$. Similarly, we pull the class y_0 defined in (12) back to a cohomology class

$$(18) \quad y''_0 \in H^1(M \setminus (K \cup K'); \mathbb{R}) \simeq H^1(D_{K \cup K'}; \mathbb{R})$$

via the inclusion $M \setminus (K \cup K') \hookrightarrow M \setminus K$. We denote

$$(19) \quad \partial_b D_{K \cup K'} = \bigcup_{\gamma \subset K_b} \Sigma_\gamma.$$

LEMMA 3.6. *If $\mu \in \mathcal{P}_\phi(D_{K \cup K'})$ satisfies $\text{supp}(\mu) \subset \partial_b D_{K \cup K'}$ then $\mu \cdot y''_0 = 0$ and $\mu \cdot y'' > 0$.*

Proof. In the construction of $D_{K \cup K'}$, as explained in Section 3.1, we fix a diffeomorphism Ψ_γ as in (1) from a small tubular neighborhood N_γ of $\gamma \subset K \cup K'$ onto $\mathbb{R}/T_\gamma\mathbb{Z} \times \mathbb{D}$, with induced tubular polar coordinates $(t, r, \theta) \in \mathbb{R}/T_\gamma\mathbb{Z} \times [0, 1] \times \mathbb{R}/2\pi\mathbb{Z}$. These coordinates model a smooth neighborhood in $D_{K \cup K'}$ of the torus component $\Sigma_\gamma = \{r = 0\} \subset \partial D_{K \cup K'}$. If $\gamma \subset K_b$ and $0 < \epsilon \ll 1$ then, since $K_b \cap K' = \emptyset$, the loop

$$t \in \mathbb{R}/T_\gamma\mathbb{Z} \mapsto (t, \epsilon, 0) \in D_{K \cup K'} \setminus \partial D_{K \cup K'} = M \setminus (K \cup K')$$

is homologous to γ in $M \setminus K'$, and the loop $\theta \in \mathbb{R}/2\pi\mathbb{Z} \mapsto (0, \epsilon, \theta)$ is homologous to zero in $M \setminus K'$. Hence, if we write $y'' \equiv pdt + qd\theta$ in $N_\gamma \setminus \gamma$, then

$$p = \frac{\langle y'', [t \mapsto (t, \epsilon, 0)] \rangle}{T_\gamma} = \frac{\langle y', \gamma \rangle}{T_\gamma} > 0$$

and

$$q = \frac{\langle y'', [\theta \mapsto (0, \epsilon, \theta)] \rangle}{2\pi} = \frac{\langle y', [\theta \mapsto (0, \epsilon, \theta)] \rangle}{2\pi} = 0.$$

Identifying $H^1(D_{K \cup K'}; \mathbb{R}) \simeq H^1(M \setminus (K \cup K'); \mathbb{R})$, we can apply Lemma 3.4 to get

$$\mu \cdot y'' = \rho^{y''}(\gamma) = \frac{\langle y', \gamma \rangle}{2\pi} > 0.$$

We now prove that $\mu \cdot y''_0 = 0$. Fix $\ell \in R_{\mathcal{F}}$ and its dual $\ell^* \in H^1(M \setminus K; \mathbb{R})$. Pulling ℓ^* back to $H^1(M \setminus (K \cup K'); \mathbb{R})$ via the inclusion $M \setminus (K \cup K') \hookrightarrow M \setminus K$, we get a class $y''_\ell \in H^1(M \setminus (K \cup K'); \mathbb{R}) \simeq H^1(D_{K \cup K'}; \mathbb{R})$. In $N_\gamma \setminus \gamma$, with $\gamma \subset K_b$, we have $y''_\ell \equiv pdt + qd\theta \equiv \ell^*$ for constants $p, q \in \mathbb{R}$, as above. If ℓ does not contain γ in its boundary then $p = q = 0$, and $\mu \cdot y''_\ell = 0$. If ℓ contains γ in its boundary then $\mu \cdot y''_\ell = \rho^{\ell^*}(\gamma) = 0$ by (II.b) in Definition 2.7. Here Lemma 3.4 was used. Since $y''_0 = \sum_{\ell \in R_{\mathcal{F}}} y''_\ell$ the desired conclusion follows. \square

3.6. Obtaining the Birkhoff section. The space $C^0(D_{K \cup K'})$ becomes a Banach space with the sup norm. Its topological dual $C^0(D_{K \cup K'})'$ is a Banach space with the corresponding dual norm, and its unit ball is a compact metrizable space when equipped with the weak* topology; here it was used $C^0(D_{K \cup K'})$ is a separable Banach space. The set $\mathcal{P}_\phi(D_{K \cup K'})$ can be seen as a convex subset of the unit ball of $C^0(D_{K \cup K'})'$, and the Riesz representation theorem implies that $\mathcal{P}_\phi(D_{K \cup K'})$ is closed in weak*. Hence $\mathcal{P}_\phi(D_{K \cup K'})$ with the weak* topology is a compact metric space.

Consider $\mathcal{E} \subset \mathcal{P}_\phi(D_{K \cup K'})$ the subset of ergodic measures. Let \mathcal{E}_b be the set of those $\mu \in \mathcal{E}$ satisfying $\mu(\partial_b D_{K \cup K'}) = 1$. Here $\partial_b D_{K \cup K'}$ is the invariant set (19). Then $\mathcal{E}_1 = \mathcal{E} \setminus \mathcal{E}_b$ is the set of those $\mu \in \mathcal{E}$ satisfying $\mu(\partial_b D_{K \cup K'}) = 0$. In general, the sets $\mathcal{E}, \mathcal{E}_b, \mathcal{E}_1$ might be non-compact in weak*.

Every cohomology class $\alpha \in H^1(D_{K \cup K'}; \mathbb{R})$ defines a function

$$\mu \in \mathcal{P}_\phi(D_{K \cup K'}) \mapsto \mu \cdot \alpha \in \mathbb{R}$$

which is weak* continuous. This is so since, by definition of intersection numbers, $\mu \cdot \alpha$ is an integral of some function in $C^0(D_{K \cup K'})$ with respect to μ . Consider the classes y'' and y_0'' defined in (17) and in (18), respectively.

Let $\mu \in \mathcal{E}_1$. By ergodicity, either $\mu(M \setminus (K \cup K')) = 1$ or $\mu(M \setminus (K \cup K')) = 0$. In the former case, μ restricts to $M \setminus (K \cup K')$ as an ergodic measure for the flow ϕ^t on $M \setminus (K \cup K')$, and Lemma 3.5 implies that $\mu \cdot y_0'' > 0$. In the latter case, we find a component $\gamma \subset K_r \cup K'$ such that $\text{supp}(\mu)$ is contained on the boundary torus Σ_γ . If $\gamma \subset K_r$ then Lemma 3.4 implies $\mu \cdot y_0'' > 0$. Here (II.a) from Definition 2.7 is also used. If $\gamma \subset K' \setminus K_r$ then we get $\mu \cdot y_0'' > 0$ from the fact that γ is a periodic orbit in the complement of K . Summarizing, we proved that

$$(20) \quad \mu \in \mathcal{E}_1 \quad \Rightarrow \quad \mu \cdot y_0'' > 0.$$

By compactness of $\mathcal{P}_\phi(D_{K \cup K'})$ we get

$$c = \sup_{\mu \in \mathcal{P}_\phi(D_{K \cup K'})} |\mu \cdot y''| < +\infty.$$

Consider

$$(21) \quad Z = \overline{\mathcal{E}_1} \cap \{\mu \in \mathcal{P}_\phi(D_{K \cup K'}) \mid \mu \cdot y_0'' = 0\}$$

where $\overline{\mathcal{E}_1}$ denotes the weak* closure. Note that Z is compact in weak* since it is a closed subset of the compact set $\mathcal{P}_\phi(D_{K \cup K'})$.

LEMMA 3.7. *If $\mu \in Z$ then $\mu \cdot y'' > 0$.*

Proof. Fix $\mu \in Z$. By the ergodic decomposition theorem we can find a Borel probability measure P_μ on \mathcal{E} such that

$$(22) \quad \mu \cdot \alpha = \int_{\mathcal{E}} \nu \cdot \alpha \, dP_\mu(\nu)$$

holds for every $\alpha \in H^1(D_{K \cup K'}; \mathbb{R})$. Apply this formula to y_0'' in combination with the definition of Z and with Lemma 3.6 to get

$$\begin{aligned} 0 = \mu \cdot y_0'' &= \int_{\mathcal{E}} \nu \cdot y_0'' \, dP_\mu(\nu) = \int_{\mathcal{E}_b} \nu \cdot y_0'' \, dP_\mu(\nu) + \int_{\mathcal{E}_1} \nu \cdot y_0'' \, dP_\mu(\nu) \\ &= \int_{\mathcal{E}_1} \nu \cdot y_0'' \, dP_\mu(\nu). \end{aligned}$$

Now (20) implies that $\nu \cdot y_0'' > 0$ for all $\nu \in \mathcal{E}_1$, and we conclude that $P_\mu(\mathcal{E}_1) = 0$. Hence $1 = P_\mu(\mathcal{E}_b)$. Substituting $\alpha = y''$ in (22) we finally get

$$\mu \cdot y'' = \int_{\mathcal{E}_b} \nu \cdot y'' \, dP_\mu(\nu) > 0$$

since by Lemma 3.6 the integrand is pointwise strictly positive. \square

By compactness of Z and Lemma 3.7, we find an open neighborhood U of Z in $\mathcal{P}_\phi(D_{K \cup K'})$ such that $\mu \cdot y'' > 0$ for all $\mu \in U$. Moreover, since (20) implies that $\mu \cdot y_0'' \geq 0$ for every $\mu \in \overline{\mathcal{E}}_1$, we get $\mu \cdot y_0'' > 0$ for every μ in $Z' = \overline{\mathcal{E}}_1 \setminus U$. By the compactness of Z'

$$d = \inf_{\mu \in Z'} \mu \cdot y_0'' > 0.$$

Finally, choose $a > 0$ large enough so that $ad > c$ and compute:

- If $\mu \in \mathcal{E}_b$ then $\mu \cdot (ay_0'' + y'') = \mu \cdot y'' > 0$.
- Assume that $\mu \in \mathcal{E}_1$. If $\mu \in U$ then $\mu \cdot (ay_0'' + y'') = a\mu \cdot y_0'' + \mu \cdot y''$ and both $\mu \cdot y_0''$ and $\mu \cdot y''$ are strictly positive. If $\mu \in \mathcal{E}_1 \setminus U \subset Z'$ then $\mu \cdot (ay_0'' + y'') \geq ad - c > 0$.

This proves that the class $y_*'' = ay_0'' + y'' \in H^1(D_{K \cup K'}; \mathbb{R})$ satisfies $\mu \cdot y_*'' > 0$ for all $\mu \in \mathcal{E}$. The ergodic decomposition theorem implies that $\mu \cdot y_*'' > 0$ must also hold for all $\mu \in \mathcal{P}_\phi(D_{K \cup K'})$. Let $y_* \in H^1(M \setminus (K \cup K'); \mathbb{R})$ be the class obtained by pulling y_*'' back via the inclusion $M \setminus (K \cup K') = D_{K \cup K'} \setminus \partial D_{K \cup K'} \hookrightarrow D_{K \cup K'}$. Then $\mu \cdot y_* > 0$ holds for all $\mu \in \mathcal{P}_\phi(M \setminus (K \cup K'))$, and Lemma 3.4 implies that $\rho^{y_*}(\gamma) > 0$ for all $\gamma \subset K \cup K'$. Hence we can apply Theorem A.1 to find the desired Birkhoff section.

4. FINDING ORBITS THAT LINK THE BROKEN BINDING

Our goal here is to prove Proposition 2.11. Let λ be a contact form on a closed, connected and oriented 3-manifold M satisfying $\lambda \wedge d\lambda > 0$, with Reeb vector field X . We consider a sequence of weighted links of periodic Reeb orbits $(\{\gamma_j^n\}, \{p_j^n\})$, $1 \leq j \leq N(n)$, as in the introduction, and the associated Borel invariant probability measures

$$\mu^n = \sum_j p_j^n \frac{(\gamma_j^n)_* \text{Leb}}{T(\gamma_j^n)}.$$

In Proposition 2.11 it is assumed that $(\{\gamma_j^n\}, \{p_j^n\})$ as above can be found so that

$$(23) \quad \mu^n \rightarrow \frac{\lambda \wedge d\lambda}{\text{vol}(\lambda)}$$

in the weak* topology of $C^0(M)'$. If we denote $I^n = \gamma_1^n \cup \dots \cup \gamma_{N(n)}^n$, then there is no loss of generality to assume that

$$(24) \quad I^n \subset I^{n+1} \quad \forall n.$$

Moreover, since L has measure zero with respect to $\lambda \wedge d\lambda$, there is no loss of generality to further assume that

$$(25) \quad I^n \subset M \setminus L \quad \forall n.$$

For every open set $U \subset M$ and $p \geq 0$, $\Omega^p(U)$ is equipped with the C_{loc}^∞ -topology, and the space $C_p(U)$ of p -currents with compact support in U is defined as the topological dual of $\Omega^p(U)$ with the associated weak* topology. Boundary operators $\partial_* : C_*(U) \rightarrow C_{*-1}(U)$ are defined as adjoints of the exterior derivative. Elements in $Z_p(U) = \ker \partial_p$ are called cycles, and the elements of $B_p(U) = \text{im} \partial_{p+1} \subset Z_p(M)$ are called boundaries. The quotient space $Z_p(U)/B_p(U)$ is canonically isomorphic to $H_p(U; \mathbb{R})$.

Let us enumerate the components of L as h_1, \dots, h_m . They can also be seen as maps $h_k : \mathbb{R}/T(h_k)\mathbb{Z} \rightarrow M$ parametrized by the flow, where $T(h_k) > 0$ is the primitive period of h_k . In this way, each h_k defines a cycle in $C_1(M)$ by the

formula $\omega \in \Omega^1(M) \mapsto \int_{\mathbb{R}/T(h_k)\mathbb{Z}} h_k^* \omega$. This cycle is also denoted by h_k , with no fear of ambiguity.

LEMMA 4.1. *Let $x_1, \dots, x_m \in \mathbb{Z}$ be such that $c = \sum_{k=1}^m x_k h_k \in B_1(M)$. Choose an oriented rational Seifert surface $f : S \rightarrow M$ for c , i.e. S is an oriented compact surface, f is an immersion, and*

- $f|_{\partial S}$ defines a submersion $\partial S \rightarrow L$ covering each h_k exactly x_k times,
- $f|_{S \setminus \partial S}$ defines an embedding $S \setminus \partial S \rightarrow M \setminus L$.

Then

$$\lim_{n \rightarrow \infty} \sum_j p_j^n \frac{\text{vol}(\lambda) \text{int}(\gamma_j^n, f)}{T(\gamma_j^n)} = \int_S f^* d\lambda = \sum_{k=1}^m x_k T(h_k)$$

where $\text{int}(\gamma_j^n, f) \in \mathbb{Z}$ denotes the algebraic intersection number between γ_j^n and f .

Proof. Consider the 3-manifold M_L obtained from M and L via the blow-up construction explained in Section 3.1. It contains a special smooth domain $D_L \subset M_L$ (7) such that $D_L \setminus \partial D_L = M \setminus L$. The vector field X extends smoothly from $M \setminus L$ to D_L tangentially to ∂D_L . The extended vector field is still denoted here by X , for simplicity. It is not difficult to show that also λ extends smoothly to D_L . Unfortunately, λ does not define a contact form on D_L since $d\lambda$ vanishes on $T_y D_L$ for every $y \in \partial D_L$. The rest of this proof is a small modification of the proof of the Action-Linking Lemma [6, Lemma 1.12].

Consider the map $\pi : D_L \rightarrow M$ defined by collapsing each boundary torus (6) Σ_{h_k} to h_k . This map is smooth since in tubular polar coordinates (t, r, θ) near $\Sigma_{h_k} = \{r = 0\}$, as in (2), the map is represented by $(t, r, \theta) \mapsto (t, r e^{i\theta})$. After a C^∞ -small perturbation of f , we can assume that $f = \pi|_S$ for an embedded surface $S \subset D_L$ that intersects ∂D_L transversely; in particular, $\partial S = S \cap \partial D_L$.

Consider a compact neighborhood U of S diffeomorphic to $[-1, 1] \times S$ in such a way that $\{0\} \times S \simeq S$. Denote by z the coordinate on $[-1, 1]$. If Ω is a positive area form on S then there is no loss of generality to assume that $\Omega \wedge dz > 0$ on U . Fix any non-negative test function φ on \mathbb{R} such that $\text{supp}(\varphi) \subset [-1, 1]$ and $\int_{\mathbb{R}} \varphi = 1$. For every $\delta \in (0, 1)$ set $\varphi_\delta(x) = \delta^{-1} \varphi(\delta^{-1} x)$. Hence $\text{supp}(\varphi_\delta) \subset [-\delta, \delta]$ and $\int_{\mathbb{R}} \varphi_\delta = 1$. Note that $\varphi_\delta(z) dz$ defines a closed 1-form supported in the interior of U and, as such, it can be smoothly extended as a closed 1-form β_δ on D_L . It represents the cohomology class dual to S . Consider smooth functions $f, g : U \rightarrow \mathbb{R}$ defined by $f = i_X dz$ and $\lambda \wedge d\lambda = g \Omega \wedge dz$. Then

$$\begin{aligned} \int_{D_L} i_X \beta_\delta \lambda \wedge d\lambda &= \int_{U \simeq [-1, 1] \times S} \varphi_\delta(z) f(z, q) g(z, q) \Omega \wedge dz \\ &= \int_{U \simeq [-1, 1] \times S} \varphi_\delta(z) (f(0, q) g(0, q) + \epsilon(z, q)) \Omega \wedge dz \\ &= \int_S f(q, 0) g(q, 0) \Omega + \int_{U \simeq [-1, 1] \times S} \varphi_\delta(z) \epsilon(z, q) \Omega \wedge dz. \end{aligned}$$

One finds $c > 0$ such that $\sup_{q \in S} |\epsilon(z, q)| \leq c|z|$ holds for all $(z, q) \in [-1, 1] \times S \simeq U$. This follows from the fundamental theorem of calculus. Hence

$$\begin{aligned} &\left| \int_{D_L} i_X \beta_\delta \lambda \wedge d\lambda - \int_S f(0, q) g(0, q) \Omega \right| \\ &\leq c\delta \left(\int_S \Omega \right) \left(\int_{[-1, 1]} \varphi_\delta(z) dz \right) = c\delta \int_S \Omega. \end{aligned}$$

Since all β_δ 's are cohomologous, the integral $\int_{D_L} i_X \beta_\delta \lambda \wedge d\lambda$ does not depend on δ . This is a simple consequence of the fact that $\lambda \wedge d\lambda$ defines an invariant measure on

$M \setminus L = D_L \setminus \partial D_L$, combined with the ergodic theorem. Taking the limit as $\delta \rightarrow 0$

$$(26) \quad \int_{D_L} i_X \beta_{\delta_0} \lambda \wedge d\lambda = \int_S fg \Omega \quad \forall \delta_0 \in (0, 1).$$

The identity $d\lambda = i_X(\lambda \wedge d\lambda)$ is valid on $D_L \setminus \partial D_L = M \setminus L$, and hence is also valid on D_L with the smooth extensions of X , λ and $d\lambda$. On U one computes

$$d\lambda = i_X(\lambda \wedge d\lambda) = i_X(g \Omega \wedge dz) = fg \Omega + \nu \wedge dz$$

for some 1-form ν . Since dz vanishes tangentially to S we get from (26) that

$$\int_{D_L} i_X \beta_\delta \lambda \wedge d\lambda = \int_S d\lambda = \sum_{k=1}^m x_k T(h_k) \quad \forall \delta \in (0, 1).$$

Finally, once $\delta \in (0, 1)$ is fixed arbitrarily, the function $i_X \beta_\delta$ is bounded on $M \setminus L$ since it is continuous on the compact space $D_L \supset D_L \setminus \partial D_L = M \setminus L$. Hence

$$\begin{aligned} \int_{D_L} i_X \beta_\delta \lambda \wedge d\lambda &= \int_{M \setminus L} i_X \beta_\delta \lambda \wedge d\lambda \\ &= \text{vol}(\lambda) \lim_{n \rightarrow \infty} \int_{M \setminus L} i_X \beta_\delta d\mu^n \\ &= \text{vol}(\lambda) \lim_{n \rightarrow \infty} \sum_j p_j^n \frac{\text{int}(\gamma_j^n, f)}{T(\gamma_j^n)}. \end{aligned}$$

Here we used the assumption that $\mu^n \rightarrow (\lambda \wedge d\lambda)/\text{vol}(\lambda)$ in weak* topology of $C^0(M)'$ in combination with Lemma B.1. \square

Let $H = \text{span}\{h_1, \dots, h_m\} \subset Z_1(M)$. Then $\dim H = m$ and $\{h_1, \dots, h_m\}$ is a basis of H . The space H contains a compact convex set

$$\mathcal{C} = \left\{ \sum_{k=1}^m a_k h_k \mid a_k \geq 0 \ \forall k, \quad \sum_{k=1}^m a_k = 1 \right\}.$$

Choose a basis e_1, \dots, e_R of $H \cap B_1(M)$. Hence $[e_r] = 0$ in $H_1(M; \mathbb{R})$.

The universal coefficients theorem tells us that $H_2(M, L; \mathbb{R}) \simeq H_2(M, L; \mathbb{Q}) \otimes \mathbb{R}$. Hence, for every r we can find finitely many $y_r^s \in \mathbb{R}$ and oriented rational Seifert surfaces f_r^s as in Lemma 4.1 such that

$$e_r = \partial \sum_s y_r^s f_r^s$$

where the f_r^s are seen in $C_2(M)$.

LEMMA 4.2. *If n is large enough then $\mathcal{C} \cap B_1(M \setminus I^n) = \emptyset$.*

Proof. There is no loss of generality to assume that $\text{vol}(\lambda) = 1$. On the space H we have the norm $\|\sum_{k=1}^m a_k h_k\| = \sum_{k=1}^m |a_k|$, and on $H \cap B_1(M)$ we have the norm $\|\sum_{r=1}^R b_r e_r\|_0 = \sum_{r=1}^R |b_r|$. The norm on $H \cap B_1(M)$ obtained by restricting $\|\cdot\|$ to $H \cap B_1(M)$ is denoted by $\|\cdot\|_1$. Since these are norms on a finite-dimensional vector space, we find $\kappa > 0$ such that $\|\cdot\|_0 \leq \kappa \|\cdot\|_1$. Choose $\eta > 0$ satisfying $\eta \geq \max_{r,s} |y_r^s|$. It is important to note that both κ and η are independent of n . By Lemma 4.1

$$n \gg 1 \quad \Rightarrow \quad \left| \sum_j p_j^n \frac{\text{int}(\gamma_j^n, f_r^s)}{T(\gamma_j^n)} - \langle f_r^s, d\lambda \rangle \right| < \frac{1}{2\kappa\eta} \min_k T(h_k) \quad \forall r.$$

If n is large enough and $c \in \mathcal{C} \cap B_1(M)$, write $c = \sum_{k=1}^m a_k h_k = \sum_{r=1}^R b_r e_r$ with $\sum_{k=1}^m a_k = 1$ and estimate

$$\begin{aligned} \left| \sum_{j,r,s} b_r y_r^s p_j^n \frac{\text{int}(\gamma_j^n, f_r^s)}{T(\gamma_j^n)} - \sum_{k=1}^m a_k T(h_k) \right| &= \left| \sum_{r,s} b_r y_r^s \left(\sum_j p_j^n \frac{\text{int}(\gamma_j^n, f_r^s)}{T(\gamma_j^n)} - \langle f_r^s, d\lambda \rangle \right) \right| \\ &\leq \left(\frac{1}{2\kappa\eta} \min_k T(h_k) \max_{r,s} |y_r^s| \right) \sum_{r=1}^R |b_r| \leq \left(\frac{1}{2\kappa} \min_k T(h_k) \right) \|c\|_0 \\ &\leq \left(\frac{1}{2} \min_k T(h_k) \right) \|c\|_1 = \left(\frac{1}{2} \min_k T(h_k) \right) \sum_{k=1}^m a_k = \frac{1}{2} \min_k T(h_k). \end{aligned}$$

Combining with

$$\sum_{k=1}^m a_k T(h_k) \geq (\min_k T(h_k)) \sum_{k=1}^m a_k = \min_k T(h_k)$$

and using the triangle inequality, we conclude that

$$(27) \quad n \gg 1 \quad \Rightarrow \quad \sum_{j,r,s} b_r y_r^s p_j^n \frac{\text{int}(\gamma_j^n, f_r^s)}{T(\gamma_j^n)} \geq \frac{1}{2} \min_k T(h_k).$$

Let Θ_n be a closed 2-form on M representing the Poincaré dual in $H^2(M; \mathbb{R})$ of the class in $H_1(M; \mathbb{R})$ represented by the cycle $\sum_j p_j^n T(\gamma_j^n)^{-1} \gamma_j^n$. Observe that for any fixed closed 1-form α on M we can compute

$$\begin{aligned} \lim_{n \rightarrow \infty} \int_M \alpha \wedge \Theta_n &= \lim_{n \rightarrow \infty} \sum_j \frac{p_j^n}{T(\gamma_j^n)} \int_{\gamma_j^n} \alpha \\ &= \lim_{n \rightarrow \infty} \int_M i_X \alpha \, d\mu^n \\ &= \frac{1}{\text{vol}(\lambda)} \int_M i_X \alpha \, \lambda \wedge d\lambda \\ &= \frac{1}{\text{vol}(\lambda)} \int_M \alpha \wedge d\lambda = 0. \end{aligned}$$

In particular, this can be applied to the closed 1-form α representing the Poincaré dual of the class in $H_2(M; \mathbb{R})$ induced by an arbitrary $S \in Z_2(M)$, thus giving

$$\lim_{n \rightarrow \infty} \langle S, \Theta_n \rangle = 0.$$

To every $A \in C_2(M)$ satisfying $\partial A \in H$ we can linearly associate \hat{A} defined by

$$\partial A = \sum_r b_r e_r \quad \Rightarrow \quad \hat{A} = \sum_{r,s} b_r y_r^s f_r^s.$$

Here we used $\partial A \in H \cap B_1(M)$. In particular, $\partial \hat{A} = \sum_r b_r e_r$ and $\partial(A - \hat{A}) = 0$.

To prove the lemma we now argue by contradiction, and assume that we can find $n_j \rightarrow \infty$ and $c_j \in \mathcal{C} \cap B_1(M \setminus I^{n_j})$. Let $A_j \in C_2(M \setminus I^{n_j})$ satisfy $\partial A_j = c_j$. By what was proved above, we can select further subsequences and assume, with no loss of generality, that

$$\left| \langle A_j - \hat{A}_j, \Theta_{n_{j+1}} \rangle \right| < \frac{1}{4} \min_k T(h_k) \quad \langle \hat{A}_j, \Theta_{n_{j+1}} \rangle \geq \frac{1}{2} \min_k T(h_k).$$

It follows that

$$(28) \quad \langle A_j, \Theta_{n_{j+1}} \rangle > \frac{1}{4} \min_k T(h_k) \quad \forall j.$$

Each A_j induces a class $[A_j] \in H_2(M, L; \mathbb{R})$. Since $H_2(M, L; \mathbb{R})$ is finite-dimensional, we find j_* such that all $[A_j]$ are in the span of $\{[A_1], \dots, [A_{j_*}]\}$. For every $j > j_*$ let us write $[A_j] = d_{j1}[A_1] + \dots + d_{jj_*}[A_{j_*}]$. Note that if $m \geq l$ then A_m is supported in $M \setminus I^{n_m}$, and by (24) we have $M \setminus I^{n_m} \subset M \setminus I^{n_l} \Rightarrow \langle A_m, \Theta_{n_l} \rangle = 0$. Hence

$$0 = \langle A_j, \Theta_{n_2} \rangle = d_{j1} \langle A_1, \Theta_{n_2} \rangle \Rightarrow d_{j1} = 0.$$

It follows that $0 = \langle A_j, \Theta_{n_3} \rangle = d_{j2} \langle A_2, \Theta_{n_3} \rangle \Rightarrow d_{j2} = 0$, and we can proceed inductively to show that $d_{js} = 0$ for every $1 \leq s \leq j_*$. This proves that $[A_j]$ vanishes in $H_2(M, L; \mathbb{R})$ when j is large enough, in contradiction to $\langle A_j, \Theta_{n_{j+1}} \rangle > 0$. \square

To conclude the proof of Proposition 2.11, fix n large so that $\mathcal{C} \cap B_1(M \setminus I^n) = \emptyset$. This is possible by the previous lemma. We invoke the Hahn-Banach theorem to obtain a continuous linear functional $\varphi_n : C_1(M \setminus I^n) \rightarrow \mathbb{R}$ satisfying

$$(29) \quad \varphi_n|_{\mathcal{C}} > 0 \quad \text{and} \quad \varphi_n|_{B_1(M \setminus I^n)} \equiv 0.$$

By reflexivity [16, §17], we find $\omega_n \in \Omega^1(M \setminus I^n)$ such that $\varphi_n(\cdot) = \langle \cdot, \omega_n \rangle$. Since φ_n vanishes on $B_1(M \setminus I^n)$ we get $d\omega_n = 0$. Since all h_k belong to \mathcal{C} , we get from (29) that

$$\int_{h_k} \omega_n > 0 \quad \forall k = 1, \dots, m.$$

The cohomology class of ω_n in $H^1(M \setminus I^n; \mathbb{R})$ is the class we were looking for, and the proof of Proposition 2.11 is complete.

5. OPENNESS OF THE SUPPORTING CONDITION

PROPOSITION 5.1. *Let X be a smooth vector field on a smooth closed 3-manifold M that has a ∂ -strong Birkhoff section $\iota : S \rightarrow M$, such that $\iota(\partial S)$ consists of non-degenerate periodic orbits. There exists a C^1 -neighborhood $\mathcal{N}(X)$ of X such that every smooth vector field $X' \in \mathcal{N}(X)$ has a ∂ -strong Birkhoff section isotopic to ι .*

Proof. Denote $K = \iota(\partial S)$, $\ell = \iota(S \setminus \partial S)$, and let $\ell^* \in H^1(M \setminus K; \mathbb{R})$ be the dual class. As in Section 3.1, we blow M up along K to get the compact manifold with boundary D_K . There is a smooth projection $P : D_K \rightarrow M$ represented as $(t, r, \theta) \mapsto (t, re^{i\theta})$ with the aid of appropriate polar tubular coordinates (2). The map P defines a diffeomorphism $D_K \setminus \partial D_K \rightarrow M \setminus K$ and collapses Σ_γ to γ . The vector field $X|_{M \setminus K}$ extends smoothly to D_K as a nonsingular vector field tangent to ∂D_K . The extended vector field is denoted by \tilde{X} . In fact, there is a boundary torus $\Sigma_\gamma \subset \partial D_K$ associated to every orbit $\gamma \subset K$, and the tubular polar coordinates (t, θ) are global periodic coordinates on $\Sigma_\gamma \simeq \mathbb{P}^+ \gamma$ that represent the projection to γ as $(t, \theta) \mapsto t$. In these coordinates \tilde{X} assumes the form $\partial_t + b(t, \theta)\partial_\theta$ and its dynamics is precisely linearized dynamics on $\mathbb{P}^+ \gamma$. It is straightforward to construct a smooth map $\tilde{\iota} : S \rightarrow D_K$ satisfying $\iota = P \circ \tilde{\iota}$. The assumption that ι is ∂ -strong tells us that $\tilde{\iota}(S)$ is an embedded surface in D_K transverse to ∂D_K . Any vector field X' on M that coincides with X on K extends smoothly to a vector field \tilde{X}' on D_K , tangent to ∂D_K . Moreover, if X' is C^1 -close to X then \tilde{X}' is C^0 -close to \tilde{X} . It follows that $\tilde{\iota}(S)$ is transverse to \tilde{X}' up to ∂D_K and there is a smooth (hence bounded) return time to $\tilde{\iota}(S)$ for the dynamics of \tilde{X}' . It follows that ι is a ∂ -strong Birkhoff section for X' provided that X' is C^1 -close enough to X .

Now take a transverse disk section D_γ near a component γ of K , where γ appears as an isolated fixed point of the first return map. By the implicit function theorem, this fixed point varies slightly under C^1 -small perturbations of X and defines a periodic orbit γ' for the perturbed vector field, C^1 -close to γ . Combined with the “extension of isotopies” theorem, we obtain that for every $\epsilon > 0$ there exists a C^1 -small neighborhood $\mathcal{N}_\epsilon(X)$ of X such that for every smooth vector field $X' \in \mathcal{N}_\epsilon(X)$

there is an ϵ -small (in C^1) smooth isotopy $(\phi_t)_{t \in [0,1]}$, $\phi_0 = id_M$, of M taking γ to a periodic orbit $\gamma' = \phi_1(\gamma)$ of X' , and similarly for the other binding components. This isotopy can be taken supported near K . By further multiplying by a function h ϵ -close to 1 in C^1 , the vector fields $(\phi_1)_*X$ and hX' get to agree on $\phi_1(K)$. Our work shows that ι is a ∂ -strong Birkhoff section for $(\phi_1)^*(hX')$. It follows that $\phi_1 \circ \iota$ is a ∂ -strong Birkhoff section for X' when $X' \in \mathcal{N}_\epsilon(X)$. \square

6. GENERICITY OF ENTROPY

Let $\phi^t : M \rightarrow M$ be a flow on a closed 3-manifold. The topological entropy $h_{top}(\phi^t)$ is a non-negative number that measures the complexity of the flow. We review a definition, originally introduced by Bowen [9]. We first endow M with a metric d . Given $T > 0$ we define $d_T(x, y) = \max\{d(\phi^t(x), \phi^t(y)) : t \in [0, T]\}$ for any couple of points $x, y \in M$. A subset $S \subset M$ is (T, ϵ) -separated if for all $x \neq y$ in S we have that $d_T(x, y) > \epsilon$. Let $N(T, \epsilon)$ be the maximal cardinality of a (T, ϵ) -separated subset of M . The topological entropy is defined as

$$h_{top}(\phi^t) = \lim_{\epsilon \rightarrow 0} \limsup_{T \rightarrow \infty} \frac{1}{T} \log(N(T, \epsilon)).$$

If M has dimension 3 and the flow is at least C^2 , deep results of Katok [35] adapted to the case of flows by Lima and Sarig [40], imply that having positive entropy is equivalent to the existence of a transverse homoclinic connection: an intersection between the stable and unstable manifolds of a hyperbolic periodic orbit that is transverse. Thus, having positive topological entropy is an open condition. Consequently, to establish Theorem 1.5 it remains to show that every contact form on a fixed closed contact 3-manifold can be arbitrarily well C^∞ -approximated by a contact form that defines the same contact structure, and whose Reeb flow has positive topological entropy.

6.1. Input from topological surface dynamics. Here we review results on the dynamics of surface homeomorphisms due to Mather [41], Koropecki [36], Koropecki, Le Calvez and Nassiri [38], and Le Calvez and Sambarino [39], as well as some basic notions in Carathéodory's theory of prime ends.

In this section S is an orientable surface without boundary. A boundary representative of S is a sequence $P_1 \supset P_2 \supset P_3 \dots$ of open sets such that ∂P_i is compact for every i , and $\cap_i P_i = \emptyset$. Here ∂P_i denotes the topological boundary of P_i relative to S . Two boundary representatives $\{P_i\}$ and $\{P'_i\}$ are equivalent if for every n there exists m such that $P_m \subset P'_n$, and *vice versa*. An ideal boundary point is an equivalence class of boundary representatives. The set of ideal boundary points of S , also called the ideal boundary of S , is denoted by $b_I S$. The ideal completion of S is defined by $c_I S = S \sqcup b_I S$, and is topologized by declaring that open subsets of S are open in $c_I S$, and that a set containing $p \in b_I S$ is a neighborhood of p if it contains $\{p\} \cup P_n$ for some n , where $\{P_i\}$ is a boundary representative of p . If $f : S \rightarrow S$ is a homeomorphism then there is an induced homeomorphism $f_S : c_I S \rightarrow c_I S$. If $b_I S$ is finite then $c_I S$ is an orientable compact surface without boundary.

Let $U \subset S$ be open. The impression of $p \in b_I U$ in S is the set

$$Z(p) = \bigcap_{\substack{V \subset c_I U \text{ open} \\ p \in V}} \text{cl}_S(V \cap U)$$

where cl_S denotes closure relative to S . Note that $Z(p)$ is closed in S and contained in the topological boundary ∂U of U relative to S . If S has finite genus then an ideal boundary point $p \in b_I U$ is said to be regular if p is isolated in $b_I U$ and $Z(p)$ has more than one point.

A compact set $K \subset S$ is called a continuum if it is connected and has at least two points. A continuum K is said to be cellular if $K = \bigcap_{n \in \mathbb{N}} D_n$ where each $D_n \subset S$ is a closed disk, and D_{n+1} is contained in the interior of D_n , for every n . Cellular continua are contractible in S . A continuum K is said to be annular if $K = \bigcap_{n \in \mathbb{N}} A_n$, where each $A_n \subset S$ is homeomorphic to a closed annulus, A_{n+1} is contained in the interior of A_n , and K separates both boundary components of A_n , for every n .

From now on S is assumed to have finite genus. Let $U \subset S$ be open and $p \in b_I U$ be regular. In Carathéodory's theory of prime ends one constructs a space $b_{\mathcal{E}_p} U$ homeomorphic to a circle, in such a way that if U is invariant by an orientation preserving homeomorphism $f : S \rightarrow S$ then there is an induced orientation preserving homeomorphism on $b_{\mathcal{E}_p} U$. Its Poincaré's rotation number is denoted by $\rho(f, p) \in \mathbb{R}/\mathbb{Z}$. The reader is referred to [38, Section 3] for a nice introduction to the theory of prime ends.

From now on let $f : S \rightarrow S$ be an orientation preserving homeomorphism. A 1-translation arc γ for f is defined to be an embedded compact arc from a point x , which is not a fixed point of f , to $f(x)$ such that $f(\gamma) \cap \gamma$ is either equal to $\{f(x)\}$ or equal to $\{x, f(x)\}$, the latter case happening precisely when x is fixed by f^2 . If $N \geq 2$ then a N -translation arc for f is a 1-translation arc such that $f^k(\gamma) \cap \gamma = \emptyset$ for all $2 \leq k \leq N-1$, and $f^N(\gamma) \cap \gamma$ is either equal to \emptyset or to $\{x\}$, the latter case happening precisely when x is fixed by f^{N+1} .

THEOREM 6.1 (Special case of Theorem 4.2 from [38]). *Suppose that S is compact, and that $U \subset S$ is an open f -invariant disk such that $S \setminus U$ has more than one point. Assume that f preserves a Borel measure μ on S such that $\mu(W) > 0$ for every non-empty open set $W \subset S$, and that $\mu(U \setminus C) < \infty$ for some compact set $C \subset U$. Let p be the point in $b_I U$ and assume that $\rho(f, p) \neq 0$. Then there exists $N \in \mathbb{N}$ and a compact set $K \subset U$ such that every N -translation arc in $S \setminus K$ is disjoint from ∂U .*

THEOREM 6.2 (Corollary 7.2 from [38]). *Suppose that S is compact, f is non-wandering, $U \subset S$ is an f -invariant open connected set, and $p \in b_I U$ is regular and fixed. If $\rho(f, p)$ is irrational then either $Z(p)$ is an annular continuum with no periodic points, or $Z(p)$ is a cellular continuum with a unique fixed point and no other periodic points.*

To conclude this section we review the notions of Moser stable and Mather sectorial periodic points from [41, Section 5]. Let q be a fixed point of f .

Choose a contracting homeomorphism α of $[0, +\infty)$, i.e. $\alpha^n(t) \rightarrow 0$ for all $t \geq 0$. The map $(s, t) \mapsto (\alpha(s), \alpha^{-1}(t))$ defines a homeomorphism of $[0, +\infty) \times [0, +\infty)$ denoted by $\alpha \times \alpha^{-1}$. An elementary sector for (f, q) is a closed subset $U \subset S$ such that $q \in \partial U$, q has a neighborhood N in U satisfying $f(N) \subset U$ and $f^{-1}(N) \subset U$, and the germ of $f|_U$ at q is topologically conjugate to the germ of $\alpha \times \alpha^{-1}$ at $(0, 0)$. Then q is said to be a Mather sectorial fixed point of f if a neighborhood of q in S is a finite union of elementary sectors for (f, q) . It turns out that this definition is independent of the choice of α .

The point q is a Moser stable fixed point of f if every neighborhood of q in S contains an f -invariant disk D with q in its interior, such that $f|_{\partial D}$ has an orbit dense in ∂D .

If q is a periodic point of f , then q is a Mather sectorial periodic point of f if it is a Mather sectorial fixed point of f^n for some n . Similarly, q is a Moser stable periodic point of f if it is a Moser stable fixed point of f^n for some n .

REMARK 6.3. Let q be a Mather sectorial periodic point of f . It follows from the above definitions that there exists $n \geq 1$ such that $f^n(q) = q$, and such that for

every $N \geq 1$ there is neighborhood V_N of q in S with the following property: every point in $V_N \setminus \{q\}$ belongs to an N -translation arc for f^n .

We continue to follow [41, Section 5] closely. A fixed connection of f is an f -invariant arc $\gamma \subset S$ whose end points are fixed points of f , or an f -invariant circle γ such that $f|_\gamma$ is orientation preserving and γ contains a fixed point of f . A periodic connection is a fixed connection of f^n , for some n .

Let μ be a Borel probability measure on S . We denote by $\mathcal{A}(S, \mu)$ the set of orientation preserving homeomorphisms $f : S \rightarrow S$ satisfying $f_*\mu = \mu$. We denote by $\mathcal{G}(S, \mu)$ the subset of $\mathcal{A}(S, \mu)$ consisting of those homeomorphisms such that every periodic point is Mather sectorial or Moser stable, and have no periodic connections.

6.2. Applications to return maps of Birkhoff sections. Our goal here is to establish the following result.

PROPOSITION 6.4. *Let λ be a strongly nondegenerate contact form defined on a closed and connected 3-manifold M . If its Reeb flow has no elliptic periodic orbits, and has a ∂ -strong Birkhoff section, then some periodic orbit has a homoclinic connection.*

We are now concerned with the proof of Proposition 6.4. The Birkhoff section is an immersion $\iota : S \rightarrow M$ defined on a compact orientable surface with boundary S as in Definition 2.4. As explained in Section 3, we can blow the link $L = \iota(\partial S)$ up and obtain a 3-manifold D_L with boundary, whose boundary components consist of invariant tori. It is straightforward to check that there is a unique immersion $\hat{\iota} : S \rightarrow D_L$ that agrees with ι on $\dot{S} = S \setminus \partial S$, maps ∂S to ∂D_L and is transverse to ∂D_L along ∂S . The proof of Proposition 5.1 shows that $\hat{\iota}$ defines a smooth embedding $S \hookrightarrow D_L$ transverse to ∂D_L that defines a global section for the extended flow on D_L , and is still denoted by S for simplicity. The return map on S is a smooth diffeomorphism. The space $c_I \dot{S}$ is nothing but the closed orientable surface obtained from S by collapsing the boundary components to points. The return map on S induces an orientation preserving homeomorphism

$$f : c_I \dot{S} \rightarrow c_I \dot{S}$$

that is smooth on \dot{S} , and such that $b_I \dot{S}$ consists of periodic points. If λ is the contact form defining the Reeb flow on M , then $d\lambda$ defines a finite Borel measure on \dot{S} that is invariant by $f|_{\dot{S}}$. After normalizing, we get an induced Borel probability measure μ on $c_I \dot{S}$ with the following properties:

- μ is f -invariant.
- Every non-empty open subset $W \subset c_I \dot{S}$ satisfies $\mu(W) > 0$.
- μ agrees on \dot{S} with a constant multiple of the measure induced by $d\lambda$.

Since every component γ of L is a hyperbolic periodic orbit of the Reeb flow, the embedded surface $S \subset D_L$ could be modified in such a way that all points in $b_I \dot{S}$ are Mather sectorial periodic points of f , keeping all the other properties described so far. This follows from a normal form of the Reeb flow near the hyperbolic closed Reeb orbit γ . Clearly, every periodic point in \dot{S} coming from a periodic orbit of the Reeb flow in $M \setminus L$ is Mather sectorial, since all closed Reeb orbits are hyperbolic. Moreover, the strong nondegeneracy assumption implies that there are no periodic connections. Summarizing, we have that $f \in \mathcal{G}(c_I \dot{S}, \mu)$.

The next statement is a direct application of [41, Theorem 5.1] and of Theorem 6.2. Its proof is extracted from the proof of [38, Theorem 8.3].

LEMMA 6.5. *Let $U \subset c_I\dot{S}$ be an open f -invariant connected set, and let $p \in b_I U$ be regular and periodic. Then $Z(p) \subset c_I\dot{S}$ is an annular continuum with no periodic points of f .*

Proof. Up to taking powers, there is no loss of generality to assume that U is invariant and p is fixed. From [41, Theorem 5.1] we get that $\rho(f, p)$ is irrational. Theorem 6.2 implies that either $Z(p)$ is an annular continuum with no periodic points, or it is a cellular continuum that has a fixed point x of f and no other periodic points of f . Note that $Z(p)$ has empty interior since it is contained in the boundary of U . Now we reproduce the argument from [38, Theorem 8.3] with small adaptations, and assume by contradiction that the latter alternative holds.

The surface $c_I\dot{S}$ cannot be a sphere. Otherwise $c_I\dot{S} \setminus Z(p)$ would be an open f -invariant disk with irrational prime ends rotation number. The fixed point x in $Z(p)$ is Mather sectorial. Since $Z(p)$ is a continuum, we can use Remark 6.3 to find, for every N , an N -translation arc γ for f^n inside any neighborhood of x , for some n , such that γ intersects $Z(p)$. This contradicts Theorem 6.1.

It follows that the universal covering $\pi : \mathbb{R}^2 \rightarrow c_I\dot{S}$ is a plane. Choose a point $\tilde{x} \in \pi^{-1}(x)$, and a lift \tilde{f} of f such that $\tilde{f}(\tilde{x}) = \tilde{x}$. Let K be the connected component of $\pi^{-1}(Z(p))$ contained in a disk \tilde{D} having \tilde{x} in its interior, and that $\pi|_{\tilde{D}}$ is a homeomorphism onto a neighborhood of $Z(p)$. Then $\tilde{f}(K) = K$ and K has empty interior. Let S' be the one-point compactification of \mathbb{R}^2 , i.e. S' is the sphere obtained from \mathbb{R}^2 by adding a point at infinity, and let $f' : S' \rightarrow S'$ be the map induced by \tilde{f} . Then $U' = S' \setminus K$ is an open disk with $\partial U' = K$ (boundary relative to S'). If p' is the (unique) point of $b_I U'$ then $\rho(f', p') = \rho(f, p)$ is irrational. But \tilde{x} is Mather sectorial and we can argue as above, using Remark 6.3, to find, for every N , an N -translation arc γ for $(f')^n$ inside any neighborhood of \tilde{x} , for some n , such that γ intersects K . This contradicts Theorem 6.1. \square

COROLLARY 6.6 (Corollary 8.7 from [38]). *If $U \subset c_I\dot{S}$ is an open connected f -periodic set with $\#b_I U < \infty$, then the boundary of U is a finite disjoint union of aperiodic annular continua and periodic points.*

Proof. The set $b_I U$ consists of finitely many periodic points for the induced map $f_U : c_I U \rightarrow c_I U$. We find $n \geq 1$ such that U is invariant by f^n , and every point in $b_I U$ is fixed by $(f_U)^n$. Let $p \in b_I U$. If $Z(p)$ is not a point then, by Lemma 6.5, it is an aperiodic annular continuum invariant by f^n . If $Z(p)$ is a point then it is a periodic point of f . The conclusion follows since the boundary of U is $\cup_{p \in b_I U} Z(p)$. \square

COROLLARY 6.7 (Corollary 8.9 from [38]). *Let $x \in \dot{S}$ be a periodic point of f , which is necessarily hyperbolic. If x belongs to some periodic continuum $K \subset c_I\dot{S}$, then the stable and unstable manifolds of x are contained in K . Moreover, all four stable and unstable branches of x have the same closure in $c_I\dot{S}$.*

Proof. Let γ be a branch of x , stable or unstable, such that $\gamma \not\subset K$. Then γ intersects a connected component V of $c_I\dot{S} \setminus K$ such that $x \in \partial V$. The set V is periodic since f preserves μ . By Corollary 6.6, ∂V is a disjoint union of periodic points and aperiodic annular continua. But the connected component of ∂V containing x is not an aperiodic continuum since it contains the periodic point x . Hence x is an isolated point of ∂V , in particular also of K , in contradiction to the fact that K is a continuum.

Let γ_1, γ_2 be branches of x . The closure K of γ_1 in $c_I\dot{S}$ is a periodic continuum. Hence, by what was proved above, $\gamma_2 \subset K$. It follows that the closure of γ_2 is contained in the closure of γ_1 . The same argument interchanging the roles of γ_1 and γ_2 concludes the proof. \square

With the above results in place, we can consider a relation on the set of periodic points of f contained in \dot{S} . Note that, by assumption, all such periodic points are hyperbolic. Similarly to [39, Section 3], two such periodic points x_1, x_2 are related if each of the four stable/unstable branches of x_1 has the same closure as any of the four branches of x_2 . This relation is clearly symmetric and transitive. Reflexivity is non-trivial and follows from Corollary 6.7. Hence, this is an equivalence relation. The set of equivalence classes is denoted by $\mathcal{E}(f)$. For a given $\kappa \in \mathcal{E}(f)$ we denote by $K(\kappa)$ the closure of one (hence any) of the four branches of a point in κ .

LEMMA 6.8 (Corollary 3.2 from [39]). *If $\kappa \in \mathcal{E}(f)$ and $x \in \dot{S}$ is a (hyperbolic) periodic point such that $x \in K(\kappa)$, then $x \in \kappa$.*

Proof. Direct consequence of Corollary 6.7. □

The above lemma was the last tool we needed to be able to conclude that the proof of [39, Proposition 5.1] can be reproduced *ipsis litteris* to establish the following statement. We do not reproduce the argument here.

LEMMA 6.9. *Let g be the genus of $c_I \dot{S}$. Every set of (necessarily hyperbolic) periodic points in \dot{S} with strictly more than $2g$ elements contains at least one point that has a homoclinic connection.*

We can now finish the proof of Proposition 6.4. Since the contact form λ is nondegenerate, a result from [10] implies that there are either two or infinitely many periodic Reeb orbits. If there are exactly two periodic orbits then the results from [12] imply that both are elliptic. Since we work under the assumption that there are no elliptic periodic Reeb orbits, it follows that there are infinitely many such periodic orbits. These give rise to infinitely many periodic points of f in \dot{S} . From Lemma 6.9 we get the desired homoclinic connection. In fact, this lemma even implies that there are infinitely many homoclinic connections.

6.3. Proof of Theorem 1.5. Fix a co-orientable contact structure on M . Consider the set Λ of contact forms defining ξ , equipped with the C^∞ -topology. This set is an open subset of the vector space $\Omega^1(M)$ equipped with the C^∞ -topology. The latter is a topological vector space whose topology can be defined by a complete metric. Consider the set $\Lambda_+ \subset \Lambda$ of contact forms whose Reeb flows have positive topological entropy. As mentioned before, Λ_+ is open, and we need to show that Λ_+ is dense in Λ .

Let $\Lambda_* \subset \Lambda$ denote the set of nondegenerate contact forms. As is well-known, Λ_* contains a countable intersection of open and dense subsets of Λ . In particular, Λ_* is dense in Λ . Let $\Lambda^e \subset \Lambda$ denote the subset of all contact forms that admit at least one nondegenerate and elliptic closed Reeb orbit. Then $\Lambda^h = \Lambda_* \setminus \Lambda^e$ consists precisely of those contact forms on Λ all of whose periodic Reeb orbits are hyperbolic.

Fix a contact form $\lambda \in \Lambda^e$, and let γ be an elliptic and nondegenerate periodic Reeb orbit of λ . By [22, Theorem 2], see also [49], there is a sequence of contact forms λ_n such that $\lambda_n \rightarrow \lambda$ in C^∞ and the Reeb flow of every λ_n has a hyperbolic closed Reeb orbit with a transverse homoclinic connection. In fact, these homoclinic connections can even be found inside any neighborhood of γ fixed *a priori*. In particular, $\lambda_n \in \Lambda_+$ for every n .

Now let $\lambda \in \Lambda \setminus \Lambda_e$. There is a sequence of contact forms λ_n satisfying $\lambda_n \rightarrow \lambda$ in C^∞ , such that every λ_n is strongly nondegenerate and has a ∂ -strong Birkhoff section for its Reeb flow. This is so because both properties are C^∞ -generic; the latter property holds in an open and dense set of contact forms in the C^∞ -topology by Corollary 1.2. If there exists a subsequence $n_j \rightarrow \infty$ such that $\lambda_{n_j} \in \Lambda_e$ then,

by what was observed so far, λ is a limit of contact forms in Λ_+ . Hence, it remains to deal with the case where there exists $N \in \mathbb{N}$ such that for every $n \geq N$ the contact form λ_n is strongly nondegenerate, its Reeb flow has a ∂ -strong Birkhoff section, and all periodic Reeb orbits are hyperbolic. Proposition 6.4 implies that for every $n \geq N$ the Reeb flow of λ_n has a hyperbolic periodic orbit with a homoclinic connection, which in turn implies that $\lambda_n \in \Lambda_+$. The proof of Theorem 1.5 is complete.

APPENDIX A. A RESULT IN SCHWARTZMAN-FRIED-SULLIVAN THEORY

In this appendix M is a closed, connected and oriented 3-manifold, and X is a smooth vector field on M . The flow of X is denoted by ϕ^t . Let $L \subset M$ be a link formed by periodic orbits. As before, we denote by $\mathcal{P}_\phi(M \setminus L)$ the set of ϕ -invariant Borel probability measures on $M \setminus L$. The statement below can be found as Theorem 2.3 in the extended arXiv version of [31], but its proof was not available in the literature as far as we know, except for some arguments sketched by Ghys [23], and for a version stated in terms of homology directions by Fried [20].

THEOREM A.1. *Suppose that there exists $y \in H^1(M \setminus L; \mathbb{R})$ satisfying*

- $\mu \cdot y > 0$ for all $\mu \in \mathcal{P}_\phi(M \setminus L)$.
- $\rho^y(\gamma) > 0$ for all $\gamma \subset L$.

Then there exists a ∂ -strong Birkhoff section for ϕ^t whose boundary components are contained in L .

REMARK A.2. A refinement of Theorem A.1 with conditions for global surfaces of section representing a prescribed homology class can be found in [29].

REMARK A.3. Some components of L may not be contained in the boundary of the Birkhoff section provided by Theorem A.1.

A.1. Preliminaries.

A.1.1. Blowing periodic orbits up. As proved in Section 3.1, we can blow L up and construct a new manifold without boundary M_L (5), containing a special smooth compact domain D_L (7). The domain $D_L \subset M_L$ is the closure of $M \setminus L$ in M_L . Moreover, X can be extended from $M \setminus L$ to M_L as a smooth vector field X_L . The vector field X_L is not unique, but its restriction to D_L is unique. Moreover, X_L is tangent to ∂D_L and the dynamics on ∂D_L captures the linearized dynamics along the components of L in a precise way described in Section 3.1.

For later purposes we need to fix some notation. Denote the components of L by $\gamma_1, \dots, \gamma_h$ and their primitive periods by $T_j > 0$. For each $j = 1, \dots, h$, $\Sigma_j \subset \partial D_L$ denotes the boundary torus associated to the end of $M \setminus L$ near γ_j .

A.1.2. Schwartzman cycles and structure currents. For every $p \geq 0$ one can turn $\Omega^p(M_L)$ into a topological vector space by equipping it with the C_{loc}^∞ -topology. Its topological dual is denoted by $C_p = \Omega^p(M_L)'$ and equipped with the weak* topology. As mentioned before, an element of C_p is a p -current with compact support. Denote by C'_p the topological dual of C_p equipped with its weak* topology. The map $\Omega^p(M_L) \rightarrow C'_p$ given by $\omega \mapsto \langle \cdot, \omega \rangle$ is a linear homeomorphism, in other words $\Omega^p(M_L)$ is reflexive. There is a boundary operator $\partial : C_{p+1} \rightarrow C_p$ defined as the adjoint of the exterior derivative $d : \Omega^p(M_L) \rightarrow \Omega^{p+1}(M_L)$. A current in C_p is called a cycle if it is in the kernel of $\partial : C_p \rightarrow C_{p-1}$, and is called a boundary if it is in the image of $\partial : C_{p+1} \rightarrow C_p$. The space of boundaries is denoted by B_p , the space of cycles by Z_p , and we have $H_p(M_L; \mathbb{R}) = Z_p/B_p$.

Consider the set $\mathcal{P}(M_L)$ of compactly supported finite Borel measures on M_L , and let $\mathcal{P}_{\phi_L}(D_L) \subset \mathcal{P}(M_L)$ be the subset of those which are ϕ_L -invariant probability measures supported on D_L . Any $\mu \in \mathcal{P}(M_L)$ defines a 1-current $c_\mu \in C_1$ by the formula

$$\langle c_\mu, \omega \rangle = \int_{M_L} \omega(X_L) d\mu, \quad \omega \in \Omega^1(M_L).$$

We follow Sullivan's notation and write $c_\mu = \int_{M_L} X_L d\mu$. If μ is supported on D_L then c_μ is a cycle if, and only if, μ is ϕ_L -invariant. The elements of the set

$$\mathcal{S}_X := \{c_\mu \mid \mu \in \mathcal{P}_{\phi_L}(D_L)\} \subset Z_1$$

will be called *Schwartzman cycles*. The *Dirac currents* $\delta_p \in C_1$, $p \in M_L$, are defined by $\langle \delta_p, \omega \rangle = \omega(X_L)|_p \in \mathbb{R}$. Let $\mathcal{C} \subset C_1$ denote the closed convex cone generated by $\{\delta_p \mid p \in D_L\}$. In [45] \mathcal{C} is called the *cone of structure currents in D_L* . The following lemma summarizes the analytical facts we need.

LEMMA A.4 ([29], Lemma 3.3 and Lemma 3.4). *The following hold.*

- (I) *There exists $\omega \in \Omega^1(M_L)$ such that $\langle c, \omega \rangle > 0$ for every $c \in \mathcal{C} \setminus \{0\}$.*
- (II) *If ω is as in (I) then the convex set $K = \{c \in \mathcal{C} \mid \langle c, \omega \rangle = 1\}$ is compact.*
- (III) *For every $c \in \mathcal{C}$ there exists a unique finite Borel measure on D_L such that $c = \int_{D_L} X_L d\mu$.*

A.2. Transverse foliations. Let β be a closed 1-form $\beta \in \Omega^1(M_L)$ that represents y on $M \setminus L = D_L \setminus \partial D_L$. We claim that

$$(30) \quad \langle c_\mu, \beta \rangle > 0 \quad \forall \mu \in \mathcal{P}_{\phi_L}(D_L).$$

To see this, let $\mu \in \mathcal{P}_{\phi_L}(D_L)$ be arbitrary. For every Borel set $E \subset M_L$ define

$$\mu_j(E) = \mu(E \cap \Sigma_j) \quad \dot{\mu}(E) = \mu(E \cap (D_L \setminus \partial D_L)) = \mu(E \cap (M \setminus L)).$$

Then $\dot{\mu}$ and the μ_j are ϕ_L -invariant Borel measures, and $\mu = \dot{\mu} + \sum_j \mu_j$. We have

$$(31) \quad \langle c_\mu, \beta \rangle = \int_{M_L} \beta(X_L) d\mu = \int_{M_L} \beta(X_L) d\dot{\mu} + \sum_{j=1}^h \int_{M_L} \beta(X_L) d\mu_j.$$

If $\mu_j(M_L) = \mu(\Sigma_j) > 0$ then $\mu_j/\mu_j(M_L) \in \mathcal{P}_{\phi_L}(D_L)$ is supported on Σ_j . By Lemma 3.4 and the hypotheses of Theorem A.1

$$\int_{M_L} \beta(X_L) d\mu_j = \int_{\Sigma_j} \beta(X_L) d\mu_j = \frac{2\pi}{T_j} \mu_j(M_L) \rho^y(\gamma_j) > 0.$$

If $\dot{\mu}(M_L) = \mu(M \setminus L) > 0$ then $\dot{\mu}/\dot{\mu}(M_L)$ induces an element of $\mathcal{P}_\phi(M \setminus L)$. By the hypotheses of Theorem A.1

$$\int_{M_L} \beta(X_L) d\dot{\mu} = \int_{M \setminus L} \beta(X) d\dot{\mu} > 0.$$

Thus each term in the sum (31) is non-negative, and at least one term is positive since $\mu(M_L) = \mu(D_L) = 1$. We have established (30).

By (I) and (II) in Lemma A.4 there exists $\omega \in \Omega^1(M_L)$ such that $\langle \cdot, \omega \rangle > 0$ on $\mathcal{C} \setminus \{0\}$, and $K = \{c \in \mathcal{C} \mid \langle c, \omega \rangle = 1\}$ is compact and convex in C_1 . Let $c \in \mathcal{C} \setminus \{0\}$ be a cycle. By (III) in Lemma A.4, $c = \int_{M_L} X_L d\nu$ for some positive finite Borel measure ν supported on D_L , and ν must be invariant because c is a cycle. In other words, $\mu := \nu/\nu(M_L) \in \mathcal{P}_{\phi_L}(D_L)$ and $c = \nu(M_L)c_\mu$ for a Schwartzman cycle c_μ . From (30) we conclude that $\langle c, \beta \rangle = \nu(M_L) \langle c_\mu, \beta \rangle > 0$. In particular $\mathcal{C} \cap B_1 = \{0\}$, or equivalently $K \cap B_1 = \emptyset$, and β evaluates positively on $K \cap Z_1$. By [29, Theorem A.1] we find $\eta_0 \in C'_1$ that vanishes on B_1 , is positive on K , and agrees with β on Z_1 . By reflexivity $C'_1 = \Omega^1(M_L)$ we conclude that η_0 is a 1-form,

and as such it must be closed since it vanishes on B_1 . Moreover, $\eta_0|_{M \setminus L}$ represents y since it agrees with β on Z_1 . Finally, note that

$$(32) \quad \eta_0(X_L)|_p = \langle \delta_p, \eta \rangle > 0 \quad \forall p \in D_L$$

because $\langle \cdot, \eta_0 \rangle > 0$ on $\mathcal{C} \setminus \{0\}$.

The kernel of η_0 integrates to a foliation transverse to X_L , and hence to X on $M \setminus L$, but in general this foliation might behave very badly. In the next Section we shall deal with this issue by well-known arguments.

A.3. Birkhoff sections. Note that $H^1(M_L; \mathbb{R}) \simeq H^1(M \setminus L; \mathbb{R})$ is a finite dimensional vector space. Hence, in view of (32), we can find arbitrarily C_{loc}^∞ -close to η_0 a 1-form η_1 on M_L which is closed, has rational periods, and still satisfies (32). Consider inclusions $\iota : D_L \hookrightarrow M_L$ and $\iota_j : \Sigma_j \hookrightarrow M_L$. The class $[\iota^* \eta_1] \in H^1(D_L; \mathbb{R})$ is close to y , and induces a class in $H^1(D_L; \mathbb{Q})$. Choose $m \in \mathbb{N}$ such that $\eta := m\eta_1$ has integer periods. Note that η satisfies (32), that is

$$(33) \quad \eta(X_L)|_p = \langle \delta_p, \eta \rangle > 0 \quad \forall p \in D_L.$$

Denote by $N \in \mathbb{N}$ a generator of the group of periods of η . Then, after arbitrarily choosing a point $p_0 \in M \setminus L$, we can define a map $\text{pr} : D_L \rightarrow \mathbb{R}/\mathbb{Z}$ by setting $\text{pr}(p)$ to be the integral of $N^{-1}\iota^*\eta$ along any path from p_0 to p , modulo \mathbb{Z} :

$$\text{pr}(p) = \frac{1}{N} \int_{p_0}^p \iota^* \eta \quad \text{mod } \mathbb{Z}.$$

The map $\text{pr} : D_L \rightarrow \mathbb{R}/\mathbb{Z}$ is a smooth surjective submersion in view of (33). It follows from this construction that if $c : S^1 \rightarrow M \setminus L$ is a smooth loop then

$$(34) \quad \frac{1}{N} \int_c \eta = \text{degree of } \text{pr} \circ c.$$

The preimages $\text{pr}^{-1}(x)$ are the leaves of a foliation of D_L obtained by integrating $\ker \iota^* \eta$. Each $\text{pr}^{-1}(x)$ is a compact embedded submanifold of D_L that intersects the boundary ∂D_L cleanly, since it is transverse to X_L and X_L is tangent to ∂D_L . It follows that $\text{pr}^{-1}(x) \subset D_L$ is a smooth embedded surface transverse to X_L with boundary equal to $\text{pr}^{-1}(x) \cap \partial D_L$. Each leaf $\text{pr}^{-1}(x)$ can be co-oriented by the vector field X_L , and can also be co-oriented by pulling back the canonical orientation of \mathbb{R}/\mathbb{Z} via the map pr . These co-orientations coincide by construction. Note that all trajectories in D_L will hit all leaves $\text{pr}^{-1}(x)$ in finite time, both in the future and in the past: this follows from compactness of D_L and from (33). In particular, for all $x \in \mathbb{R}/\mathbb{Z}$ the surface $\text{pr}^{-1}(x)$ is a global surface of section for the flow of X_L on D_L .

In a final step we modify one given fiber of the map pr to obtain a Birkhoff section for the flow of X on M . Let us fix $x \in \mathbb{R}/\mathbb{Z}$ arbitrarily. For each $j \in \{1, \dots, h\}$, we choose coordinates $(t, \theta) \in \mathbb{R}/T_j\mathbb{Z} \times \mathbb{R}/2\pi\mathbb{Z} \simeq \Sigma_j$ as explained in Section 3.1. These can be used to write a basis $\{dt, d\theta\}$ of $H^1(\Sigma_j; \mathbb{R})$. Note that $N^{-1}\iota_j^*\eta$ is homologous to $a_1 dt/T_j + a_2 d\theta/2\pi$ for some $a_1, a_2 \in \mathbb{Z}$. Assume that $\text{pr}^{-1}(x) \cap \Sigma_j \neq \emptyset$, and let us consider a connected component α of $\text{pr}^{-1}(x) \cap \Sigma_j$, oriented as part of the boundary of $\text{pr}^{-1}(x)$. Then α is non-trivial in $H_1(\Sigma_j; \mathbb{Z})$, since otherwise it would bound a disk $D \subset \Sigma_j$ with $X_L \pitchfork \partial D$ thus forcing a singularity of X_L on Σ_j . Let $\{e_1, e_2\}$ be the basis in $H_1(\Sigma_j; \mathbb{Z})$ dual to $\{dt/T_j, d\theta/2\pi\}$ and write $\alpha = n_1 e_1 + n_2 e_2$ in homology. We already know that $(n_1, n_2) \neq (0, 0)$. Since α is an embedded loop in Σ_j we also know that (n_1, n_2) is primitive: if $d \in \mathbb{N}$ and $(n_1, n_2)/d \in \mathbb{Z} \times \mathbb{Z}$ then $d = 1$. Moreover, $0 = N^{-1} \int_\alpha \iota_j^* \eta = a_1 n_1 + a_2 n_2$. The numbers n_1, n_2 depend only on j and not on the choice of α . Using the transversality of the flow, we can deform $\text{pr}^{-1}(x)$ near $\text{pr}^{-1}(x) \cap \Sigma_j$ to obtain a surface S that intersects ∂D_L transversely, is transverse to X_L up to the boundary, still is a global section for

the flow on D_L , and such that the following holds: if $n_1 = 0$ then S coincides with finitely many annuli of the form $\{t = t_*, r \in [0, \epsilon)\}$ near Σ_j ; if $n_1 \neq 0$ then dt does not vanish tangentially to $S \cap \Sigma_j$. Now we collapse each Σ_j to γ_j to obtain again the manifold M from D_L . This collapsing map can be defined by $(t, r, \theta) \mapsto (t, r e^{i\theta})$ near each Σ_j . In this process S gets mapped to a Birkhoff section for the flow of X on M . The transversality to the flow on ∂D_L before applying the projection implies that we get a ∂ -strong Birkhoff section. In the notation above, if $n_1 = 0$ then each loop α gets mapped to an interior point of the obtained section where it intersects γ_j transversely, and if $n_1 \neq 0$ then the orbit γ_j gets covered $|n_1|$ times.

REMARK A.5. Since the Birkhoff section obtained is ∂ -strong, the associated return time function is bounded.

APPENDIX B. A LEMMA ON WEAK* CONVERGENCE

Let X be a compact manifold and let μ be a regular¹ Borel probability measure on X . Let μ_n be a sequence of Borel probability measures satisfying $\mu_n \rightarrow \mu$ in the weak* topology.

LEMMA B.1. *If $V \subset X$ is open with $\mu(\partial V) = 0$, and if $f : V \rightarrow \mathbb{R}$ is a continuous bounded function, then*

$$\int_V f d\mu_n \rightarrow \int_V f d\mu.$$

Proof. We start with a claim.

Claim. For every $\eta > 0$ there exists an open neighborhood W of ∂V , and $n_0 \geq 1$, such that $n \geq n_0 \Rightarrow \mu_n(W) < \eta$.

To prove the claim, we use the regularity of μ to find an open neighborhood W' of ∂V such that $\mu(W') < \eta$. Now we find $\psi : X \rightarrow [0, 1]$ continuous such that $\text{supp}(\psi) \subset W'$ and $\psi|_W \equiv 1$ on some open neighborhood $W \subset W'$ of ∂V . Then

$$\eta > \mu(W') \geq \int_X \psi d\mu = \lim_{n \rightarrow \infty} \int_X \psi d\mu_n$$

which implies that we can find n_0 such that

$$n \geq n_0 \Rightarrow \mu_n(W) \leq \int_X \psi d\mu_n < \eta$$

as desired. The proof of the claim is complete.

Let $\epsilon > 0$ be fixed arbitrary. For any $\delta > 0$ we can use the regularity of μ to find a compact set $K_0 \subset V$ such that $\mu(K_0) > \mu(V) - \delta/2$, and an open set $U_0 \supset \bar{V}$ such that $\mu(U_0) < \mu(\bar{V}) + \delta/2$. By the claim proved above we find an open neighborhood W of ∂V and n_0 such that $\mu_n(W) < \delta$ for all $n \geq n_0$. Consider sets

$$K = K_0 \cup (V \setminus W), \quad U = U_0 \cap (V \cup W).$$

Then $K \supset K_0$ is compact, $U \subset U_0$ is open, and $U \setminus K \subset W$. Moreover, we can estimate

$$\mu(U \setminus K) \leq \mu(U_0 \setminus K_0) = \mu(U_0) - \mu(K_0) < \mu(\bar{V}) + \delta/2 - (\mu(V) - \delta/2) = \delta$$

where the assumption $\mu(\partial V) = 0$ was used to get the equality $\mu(V) = \mu(\bar{V})$. In addition, we have

$$n \geq n_0 \Rightarrow \mu_n(U \setminus K) \leq \mu_n(W) < \delta.$$

Now consider $\varphi : X \rightarrow [0, 1]$ continuous such that $\text{supp}(\varphi) \subset U \setminus K$ and $\varphi \equiv 1$ on some neighborhood of ∂V . Let $h : X \rightarrow \mathbb{R}$ be the continuous function that agrees

¹A Borel probability measure is regular if for every Borel set E and every $\epsilon > 0$ we find a compact set K and an open set U such that $K \subset E \subset U$, $\mu(K) > \mu(E) - \epsilon$ and $\mu(U) < \mu(E) + \epsilon$.

with $(1 - \varphi)f$ on V and vanishes on $X \setminus V$. By the assumption that $\mu_n \rightarrow \mu$ weak*, we can find n_1 such that

$$n \geq n_1 \Rightarrow \left| \int_X h d\mu_n - \int_X h d\mu \right| < \delta.$$

We can estimate

$$\begin{aligned} \left| \int_V f d\mu_n - \int_V f d\mu \right| &= \left| \int_V \varphi f d\mu_n - \int_V \varphi f d\mu + \int_X h d\mu_n - \int_X h d\mu \right| \\ &\leq \left| \int_V \varphi f d\mu_n \right| + \left| \int_V \varphi f d\mu \right| + \left| \int_X h d\mu_n - \int_X h d\mu \right| \\ &\leq (\sup |f|) (\mu_n(U \setminus K) + \mu(U \setminus K)) + \left| \int_X h d\mu_n - \int_X h d\mu \right| \end{aligned}$$

from where we see that if $n \geq \max\{n_0, n_1\}$ then

$$\left| \int_V f d\mu_n - \int_V f d\mu \right| \leq 2\delta (\sup |f|) + \delta.$$

Finally, note that we could have chosen $\delta > 0$ satisfying $2\delta (\sup |f|) + \delta < \epsilon$. The proof of the lemma is complete. \square

APPENDIX C. AN ALTERNATIVE PROOF OF THE C^1 -DENSITY OF THE EXISTENCE OF A BIRKHOFF SECTION

In this appendix we sketch an alternative proof of the following direct consequence of Corollary 1.2. As the proof takes a different path, we think it is worth mentioning it.

THEOREM C.1. *If (M, ξ) is a closed oriented 3-manifold then the set of Reeb vector fields for ξ on M that admit a Birkhoff section is C^1 -dense in the set of Reeb vector fields.*

As we have mentioned in the introduction, a C^∞ -generic contact form λ on (M, ξ) is nondegenerate. Let R be the Reeb vector field of λ . In that case, by Theorem 1.1 of [10], R is carried by a broken book decomposition (K, \mathcal{F}) . The binding K of the broken book decomposition has radial and broken components, and if all the components are radial the broken book decomposition is a rational open book decomposition and any page is a Birkhoff section. The proof of Theorem 4.13 in [10] provides a mechanism to successively eliminate broken components of the binding K , inspired by Fried [21]. It works as follows, see also Figure 2.

Let $k \in K$ be a broken component of the binding, which is automatically a hyperbolic orbit of the Reeb flow R . Assume it is positive (i.e. the eigenvalues of the linearized first return map are positive) and denote by N and S the two components of the complement of k in its stable manifold, and by E and W the two components of the complement of k in its unstable manifold (in red and blue respectively on Figure 2 left). If there is a transverse homoclinic connection p from E to N and a transverse homoclinic connection q from W to S , then one can find a periodic orbit p' of R close to p and another one q' close to q . There is also a periodic orbit r' following successively p and q thanks to the combinatorial description of the dynamics. In this situation, Fried [21] constructed a pair of pants P transverse to R , bounded by these three orbits, and intersecting the orbit k in its interior, see Figure 2.

We add this section P to our broken book decomposition, via the following process, sometimes called Fried sum: denoting by $S_{R_{\mathcal{F}}}$ the union of the rigid pages, we consider the union $P \cup S_{R_{\mathcal{F}}}$. It is an immersed surface with 1-dimensional singularities consisting of arcs ending on boundary orbits of P or $S_{R_{\mathcal{F}}}$, see Figure 3

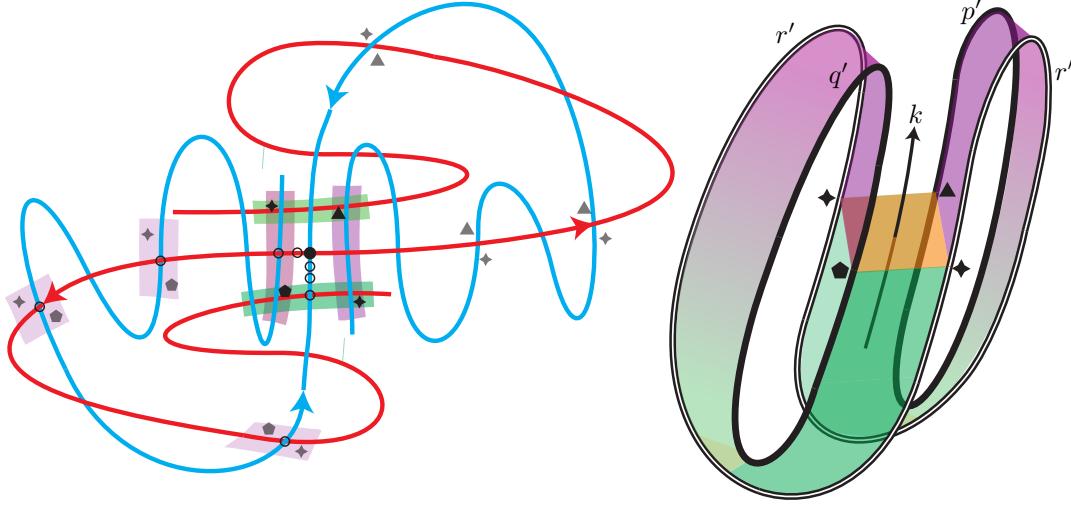


FIGURE 2. On the left, a transversal picture where the central dot is k , the stable manifolds are in blue, the unstable ones in red, the empty dots correspond to the homoclinic W-S connection q , the pentagons to the orbit q' , the triangles to the orbit p' and the 4-pointed stars to the orbit r' . On the right the pair of pants P transverse to the vector field R and intersecting the orbit k constructed by Fried: it is the union of a parallelogram in a local transverse disc whose vertices lie on p', r', q' and r' respectively, and the image under the flow of the two “almost stable edges” of this parallelogram. This topological surface is not strictly speaking transverse to the flow, but one can smooth it and make it transverse to R at the same time.

left. We desingularize those arcs transversally to the vector field, as on the central picture. On a broken binding orbit transverse to P , the result is depicted on Figure 3 right: the addition of a meridian circle corresponding to the intersection with P to the boundary of $S_{R_{\mathcal{F}}}$ gives a boundary for the new surface that intersects all orbits of the projectivized flow, including the stable and unstable direction of the broken binding orbit. We then obtain a new broken book decomposition, whose new binding is the initial one increased by the three new orbits p', q', r' and whose broken binding is the previous with k removed, i.e. k is no more a broken binding component anymore: it can either disappear from the binding or become a radial component, depending on the balance between the number of incoming and outgoing rigid pages. Since the orbits p', q' and r' are intersecting the pages of (K, \mathcal{F}) transversally, they become radial components of the new binding.

By applying this process on all broken binding orbits successively, we reduce the broken part of the binding. After a finite number of steps, the broken book has only radial binding components, and all its pages are Birkhoff sections for the Reeb flow.

The case where k is a negative hyperbolic periodic orbit is treated in the same manner. The difference is that now one needs to consider the second iterate of the return map to a local transversal to the periodic orbit in order to have the same picture.

The conclusion so far is that the proof of Theorem C.1 boils down to proving that by a C^1 -small perturbation supported away from the binding K , one can introduce

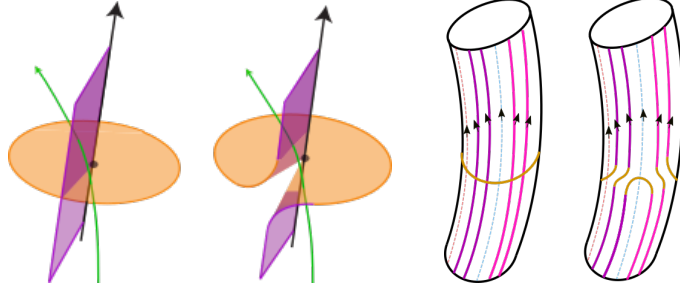


FIGURE 3. On the left is depicted the Fried sum of two transverse surfaces, which amounts to consider their union and desingularize it transversally to the vector field. On the right what happens along a broken binding orbit k when adding the pair of pants P to the collection of rigid pages $S_{R_{\mathcal{F}}}$. The torus is the blow-up of k . The stable and unstable manifolds are dotted, the rigid pages are in pink and purple depending on the sector, and P is the orange meridian circle. If the sum of the longitudinal coordinates of all boundary components of $S_{R_{\mathcal{F}}}$ add up to 0, then the Fried union with P also has zero longitudinal coordinate, hence the boundary of the new surface can be made meridional. After an isotopy, we obtain a surface with no boundary component along k and transverse to it. If this sum is nonzero, then the obtained surface still has k in its boundary, but k is now in the radial part of the boundary (*i.e.* the flow winds with respect to the new surface).

transverse homoclinic connections at will. Note that if the perturbation is small enough and supported outside a fixed neighborhood of K , then the Reeb vector field remains carried by (K, \mathcal{F}) since the transversality condition is open away from the binding.

This is a direct application of the work of Bonatti-Crovisier [8] and Arnaud-Bonatti-Crovisier [4] in the Reeb case. The only thing to check is that Reeb flows satisfy the *lift Axiom* of Pugh-Robinson [43]. For its statement and proof below, we denote by $D_r(z) \subset \mathbb{R}^{2n}$ and $B_r(z) \subset \mathbb{R}^{2n}$ the closed and open euclidean balls of radius $r > 0$ and center $z \in \mathbb{R}^{2n}$, respectively.

LEMMA C.2 (Lift Axiom). *Equip $\mathbb{R}^{2n} \times \mathbb{R}$ with coordinates (z, t) . Consider a contact form λ on $D_1(0) \times [0, 1]$ with Reeb vector field $R = \partial_t$. There exist $K > 0$ and $0 < \epsilon_* < 1/2$ with the following property. For every $0 < \epsilon < \epsilon_*$ and $z_0 \in B_\epsilon(0) \setminus \{0\}$ there exists a smooth contact form λ' on $D_1(0) \times [0, 1]$ with the same contact kernel than λ such that $\|\lambda' - \lambda\|_{C^2} < K\epsilon$, $\lambda' - \lambda$ is compactly supported in the open box $B_{\epsilon^{-1}|z_0|}(0) \times (0, 1)$, and the Reeb flow of λ' takes $(0, 0)$ to $(z_0, 1)$.*

Proof. Fix $\beta : [0, 1] \rightarrow [0, 1]$ smooth such that $\beta(t) = 0$ for t near 0, $\beta(t) = 1$ for t near 1, $\beta' \geq 0$. For every $z_0 \in B_1(0)$ consider the curve $\gamma(t; z_0) = (\beta(t)z_0, t)$. We look for a family of contact Hamiltonians $h_{z_0}(z, t)$ on $D_1(0) \times [0, 1]$, depending on a parameter $z_0 \in \mathbb{R}^{2n}$. The associated contact Hamiltonian vector field is $X = h_{z_0}R + Y$, where Y is determined by $i_Y\lambda = 0$, $dh_{z_0} - (i_R dh_{z_0})\lambda = i_Y d\lambda$. We would like to achieve that X is positively tangent to γ , and that $h_{z_0} = 1$ along γ . A simple calculation shows that these constraints can be met if $|z_0|$ is small enough. Moreover, if we represent $dh_{z_0}(\gamma(t)) = \langle V(t; z_0), \cdot \rangle$ by a vector field then $|V| + |\partial_t V| + |\partial_t^2 V| \leq c|z_0|$ for some $c > 0$ independent of z_0 . Note that $\langle V(t; z_0), \partial_t \gamma(t; z_0) \rangle = 0$ for all t . Perhaps after making c larger, but still

independent of z_0 , the function $\hat{h}_{z_0}(z, t) = 1 + \langle V(t; z_0), (z - \beta(t)z_0, 0) \rangle$ satisfies

$$\|\hat{h}_{z_0} - 1\|_{C^2(D_1(0) \times [0,1])} \leq c|z_0|, \quad |\hat{h}_{z_0}(z, t) - 1| \leq c|z_0|(|z| + |z_0|)$$

and meets the requirements along γ . The discussion so far is independent of ϵ . Consider smooth bump functions $\phi_{z_0}^\epsilon(z)$ with values in $[0, 1]$, $\phi_{z_0}^\epsilon = 1$ on $D_{2|z_0|}(0)$, $\text{supp}(\phi_{z_0}^\epsilon) \subset B_{\epsilon^{-1}|z_0|}(0)$, $\|\nabla\phi_{z_0}^\epsilon\|_\infty = O(\epsilon|z_0|^{-1})$, $\|\nabla^2\phi_{z_0}^\epsilon\|_\infty = O(\epsilon^2|z_0|^{-2})$. To conclude, note that there exists $\epsilon_* > 0$ small enough such that if $0 < \epsilon < \epsilon_*$ then $h_{z_0} = 1 + \phi_{z_0}^\epsilon(z)(\hat{h}_{z_0}(z, t) - 1)$ is the contact Hamiltonian we were looking for, and $\lambda' = h_{z_0}\lambda$. \square

For the rest of the argument, recall from [4, 8] that it uses the fact that in the conservative setting there are no wandering points and that there exists $N > 0$ depending only on the C^1 -size of the perturbation ϵ , such that given any fixed closed neighborhood U of finitely many periodic orbits one can join any two points in the same connected component of $M \setminus (\cup_{t \in [0, N]} \phi_t(U))$ by a δ -pseudo-orbit, where $\delta \ll \epsilon$ can be chosen small at will independently of N and ϵ . The jumps of the pseudo-orbit are contained in a finite number of disjoint flow-boxes in $M \setminus U$ whose radii are of order δ and length N . To construct the flow-boxes, δ has to be taken very small once ϵ and N have been fixed. Here ϕ_t denotes the flow of R at time t . The δ -pseudo-orbits can then be turned into genuine ones by ϵ - C^1 -small and δ - C^0 -small perturbations of the flow, supported in the flow-boxes (and thus in the complement of U) and given by the lift Axiom, see [4, 8].

Precisely here, to make sure that we create a homoclinic connection, we take a point a on a unstable manifold of a broken component of the binding k and a point b on one of its stable manifolds. The set U is a small closed neighborhood of the binding of the supporting broken book so that a and b are in the boundary of U . We can moreover get that

- the orbit of a for negative time stays in U as well as the orbit of b for positive times.
- the orbit of a for time $[0, N]$ stays outside of U as well as the orbit of b for time $[-N, 0]$.

We can then find, following [4, 8], two first flow-boxes around the portions of orbits $\phi_{[0, N]}(a)$ and $\phi_{[-N, 0]}(b)$ and the remaining ones away from the two first and U , provided the endpoints $\phi_N(a)$ and $\phi_{-N}(b)$ belong to the same connected component of $M \setminus (\cup_{t \in [0, N]} \phi_t(U))$. This last condition is always achievable by taking U small enough.

Then the Lift Axiom gives a perturbed Reeb vector field, with a deformation supported in our collection of flow-boxes (and thus in $M \setminus U$), whose flow connects a and b . This in turn yields a homoclinic orbit for k passing through a and b , since we did not perturb the portions of orbits in U through a and b that were respectively negatively and positively asymptotic to k . An extra C^∞ -small perturbation is enough to make the homoclinic connection transverse. If δ is taken small enough, the transversality of the perturbed Reeb vector field with the pages of the broken book away from the binding is preserved.

To go slightly further, let us summarize the connecting lemma that is in particular obtained in the previous discussion:

THEOREM C.3 (Arnaud, Bonatti and Crovisier [4] + ϵ). *Let $(M, \xi = \ker \lambda)$ be a closed contact 3-manifold with λ nondegenerate, then given any two points $x, y \in M$, there exists a C^2 -small perturbation λ' of λ such that y is in the positive orbit of x under the flow of $R_{\lambda'}$.*

From this result, following Arnaud-Bonatti-Crovisier [4, Section 5.1], one obtains:

THEOREM C.4. *On a closed contact 3 manifold, the set of transitive Reeb vector fields is a G_δ -dense in the set of Reeb vector fields for the C^1 -topology.*

Proof. The arguments of [4, Section 5.1] apply *verbatim*. We reproduce them here for the reader's convenience. First the set \mathcal{G}_0 of nondegenerate Reeb vector fields is a G_δ . We then consider a countable basis $(U_n)_{n \in \mathbb{N}}$ of neighborhoods of M and for $m, n \in \mathbb{N}$ we let $\mathcal{U}_{m,n}$ be the set of Reeb vector fields R for which the flow $(\phi_t)_{t \in \mathbb{R}}$ satisfies: there exists $t > 0$ with $\phi_t(U_n) \cap U_m \neq \emptyset$. The set $\mathcal{U}_{m,n}$ is open in the C^1 -topology. Moreover $\mathcal{G}_0 \cap \mathcal{U}_{m,n}$ is C^1 -dense by the connecting lemma (Theorem C.3) and so is $\mathcal{U}_{m,n}$. Thus

$$\mathcal{G} = \bigcap_{m,n} \mathcal{U}_{m,n}$$

is a G_δ .

Every R in \mathcal{G} is by definition topologically transitive and thus transitive: it has a dense orbit. Indeed, if $R \in \mathcal{G}$, for every $m \in \mathbb{N}$ the set V_m of points of M whose positive orbit meets U_m is open and dense. Thus $\bigcap_m V_m$ is a G_δ . Every point $q \in \bigcap_m V_m$ is on a dense orbit. \square

REFERENCES

- [1] M. R. R. Alves, V. Colin and K. Honda, *Topological entropy for Reeb vector fields in dimension three via open book decompositions*. Jour. École Polytechnique, Tome 6 (2019), 119–148. – cited on p. 4.
- [2] M. R. R. Alves and M. Meiwes, *Dynamically exotic contact spheres in dimensions ≥ 7* . Comment. Math. Helv. 94 (2019), no. 3, 569–622. – cited on p. 4.
- [3] M. R. R. Alves and A. Pirnapasov, *Reeb orbits that force topological entropy*. Ergodic Theory & Dynamical Systems (2021), 1–44. doi:10.1017/etds.2021.80 – cited on p. 4.
- [4] M.-C. Arnaud, C. Bonatti and S. Crovisier, *Dynamiques symplectiques génériques*, Ergodic Theory & Dynamical Systems 25 (2005), 1401–1436. – cited on p. 3, 34, 35, 36.
- [5] K. L. Baker, J. B. Etnyre and J. van Horn-Morris, *Cabling, contact structures and mapping class monoids*. J. Differential Geom. 90 (2012), no. 1, 1–80. – cited on p. 9.
- [6] D. Bechara-Senior, U. Hryniewicz, P. A. S. Salomão, *On the relation between action and linking*. J. Mod. Dyn. 17 (2021) 319–336. – cited on p. 5, 19.
- [7] G. D. Birkhoff, *Dynamical Systems*, AMS Collog. Publ. IX, Providence (1966). – cited on p. 2, 3.
- [8] C. Bonatti and S. Crovisier, *Réurrence et genericité*, Invent. Math. 158 (2004), 33–104. – cited on p. 3, 34, 35.
- [9] R. Bowen, *Entropy for Group Endomorphisms and Homogeneous Spaces*, Trans. Amer. Math. Soc. 153 (1971), 401–414. – cited on p. 23.
- [10] V. Colin, P. Dehornoy and A. Reichtman, *On the existence of supporting broken book decompositions for contact forms in dimension 3*. arXiv:2001.01448v3 – cited on p. 2, 3, 4, 5, 6, 7, 8, 9, 11, 27, 32.
- [11] D. Cristofaro-Gardiner, M. Hutchings and D. Pomerleano, *Torsion contact forms in three dimensions have two or infinitely many Reeb orbits*. Geom. Topol. 23 (2019), no. 7, 3601–3645. – cited on p. 2.
- [12] D. Cristofaro-Gardiner, U. Hryniewicz, M. Hutchings and H. Liu, *Contact three-manifolds with exactly two simple Reeb orbits*. arXiv:2102.04970v3 – cited on p. 27.
- [13] G. Contreras and M. Mazzucchelli, *Existence of Birkhoff sections for Kupka-Smale Reeb flows of closed contact 3-manifolds*. arXiv:2110.07491v2 – cited on p. 3, 5, 6.
- [14] P. Dehornoy, *Genus one Birkhoff sections for geodesic flows*. Ergod. Th. Dynam. Sys. 35 (2014) 1795–1813. – cited on p. 3.
- [15] P. Dehornoy and M. Shannon, *Almost equivalence of algebraic Anosov flows*. arXiv:1910.08457 – cited on p. 3.
- [16] G. de Rham, *Variétés différentiables. Formes, courants, formes harmoniques*. Paris: Hermann 1955. – cited on p. 22.
- [17] A. Fathi, F. Laudenbach and V. Poénaru, *Travaux de Thurston sur les surfaces, Séminaire Orsay*, Astérisque, 66-67, Société Mathématique de France, Paris, 1979. 284 pp. – cited on p. .
- [18] A. Florio and U. Hryniewicz, *Quantitative conditions for right-handedness*. arXiv:2106.12512v1 – cited on p. 7.

- [19] U. Frauenfelder and F. Schlenk, *Fiberwise volume growth via Lagrangian intersections*. J. Symplectic Geom., 4 (2006), 117–148. – cited on p. 4.
- [20] D. Fried, *The geometry of cross sections to flows*. Topology, 21 (1982), 353–371. – cited on p. 4, 10, 28.
- [21] D. Fried, *Transitive Anosov flows and pseudo-Anosov maps*, Topology, 22 (1983), 299–303. – cited on p. 2, 3, 5, 32.
- [22] V. Gelfreich and D. Turaev, *Universal dynamics in a neighborhood of a generic elliptic periodic point*. Regul. Chaotic Dyn. 15 (2010), no. 2-3, 159–164. – cited on p. 27.
- [23] E. Ghys, *Right-handed vector fields & the Lorenz attractor*, Japan. J. Math. 4, 47–61 (2009). – cited on p. 28.
- [24] E. Giroux, *Géométrie de contact: de la dimension trois vers les dimensions supérieures*. Proceedings of the International Congress of Mathematicians, Vol. II (Beijing, 2002), 405–414, Higher Ed. Press, Beijing, 2002. – cited on p. 1.
- [25] G. Hector and U. Hirsch, *Introduction to the geometry of foliations. Part B. Foliations of codimension one*, Second edition. Aspects of Mathematics, E3. Friedr. Vieweg & Sohn, Braunschweig, 1987. x+297 pp. – cited on p. 11.
- [26] H. Hofer, K. Wysocki and E. Zehnder. *The dynamics of strictly convex energy surfaces in \mathbb{R}^4* . Ann. of Math. 148 (1998), 197–289. – cited on p. 2.
- [27] H. Hofer, K. Wysocki and E. Zehnder, *Finite energy foliations of tight three-spheres and Hamiltonian dynamics*. Ann. Math 157 (2003), 125–255. – cited on p. 3, 4, 6, 9.
- [28] U. Hryniewicz. *Systems of global surfaces of section for dynamically convex Reeb flows on the 3-sphere*. J. Symplectic Geom. 12 (2014), 791–862. – cited on p. 3.
- [29] U. Hryniewicz, *A note on Schwartzman-Fried-Sullivan Theory, with an application*. J. Fixed Point Theory Appl. 22 (2020), no. 1, Paper No. 25, 20 pp. – cited on p. 3, 4, 7, 10, 13, 14, 28, 29.
- [30] U. Hryniewicz and P. A. S. Salomão. *On the existence of disk-like global sections for Reeb flows on the tight 3-sphere*. Duke Math. J. 160, 3 (2011), 415–465. – cited on p. 3.
- [31] U. Hryniewicz and P. A. S. Salomão. *Global surfaces of section for Reeb flow in dimension three and beyond*. Proc. Int. Cong. of Math. – 2018, Rio de Janeiro, Vol. 1 (937–964). – cited on p. 28.
- [32] U. Hryniewicz, P. A. S. Salomão and K. Wysocki, *Genus zero global surfaces of section for Reeb flows and a result of Birkhoff*. To appear in J. Eur. Math. Soc., [arXiv:1912.01078v3](https://arxiv.org/abs/1912.01078v3) – cited on p. 3.
- [33] M. Hutchings, *Lectures notes on embedded contact homology*. Contact and symplectic topology, 389–484, Bolyai Soc. Math. Stud. 26, Springer, 2014. – cited on p. 2.
- [34] K. Irie, *Equidistributed periodic orbits of C^∞ -generic three-dimensional Reeb flows*. J. Symplectic Geom. 19 (2021), no. 3, 531–566. – cited on p. 2.
- [35] A. Katok, *Lyapunov exponents, entropy and periodic orbits for diffeomorphisms*. Publications mathématiques de l’I.H.É.S., tome 51 (1980), p. 137–173. – cited on p. 23.
- [36] A. Koropecski, *Aperiodic invariant continua for surface homeomorphisms*. Math. Z. 266 (2010), no. 1, 229–236. – cited on p. 23.
- [37] G. Knieper and H. Weiss, *C^∞ genericity of positive topological entropy for geodesic flows on S^2* . J. Diff. Geom. 62 (2002), 127–141. – cited on p. 4.
- [38] A. Koropecski, P. Le Calvez and M. Nassiri, *Prime ends rotation numbers and periodic points*. Duke Math. J. 164 (2015), no. 3, 403–472. – cited on p. 4, 23, 24, 25, 26.
- [39] P. Le Calvez, M. Sambarino, *Homoclinic orbits for area preserving diffeomorphisms of surfaces*. [arXiv:1912.06792v1](https://arxiv.org/abs/1912.06792v1) – cited on p. 4, 23, 27.
- [40] Y. Lima and O. Sarig, *Symbolic dynamics for three-dimensional flows with positive topological entropy*, J. Eur. Math. Soc. (JEMS) 21 (2019), no. 1, 199 – 256. – cited on p. 23.
- [41] J. N. Mather, *Invariant subsets for area preserving homeomorphisms of surfaces*. Mathematical analysis and applications, Part B, Adv. in Math. Suppl. Stud., vol. 7, Academic Press, New York, 1981, pp. 531–562. – cited on p. 23, 24, 25, 26.
- [42] L. Polterovich, *Growth of maps, distortion in groups and symplectic geometry*. Invent. Math. 150 (2002), no. 3, 655–686. – cited on p. 4.
- [43] C. Pugh and C. Robinson, *The C^1 -closing Lemma, including Hamiltonians*, Erg. Th. Dynam. Syst. 3 (1983), 261–313. – cited on p. 34.
- [44] S. Schwartzman, *Asymptotic cycle*, Ann. of Math. (2) 66 (1957), 270–284. – cited on p. 4.
- [45] D. Sullivan, *Cycles for the dynamical study of foliated manifolds and complex manifolds*. Invent. Math. 36 (1976), 225–255. – cited on p. 4, 29.
- [46] K. Sigmund, *On the space of invariant measures for hyperbolic flows*. Amer. J. Math. 94, No. 1 (1972), 31–37. – cited on p. 3.
- [47] C. H. Taubes, *The Seiberg-Witten equations and the Weinstein conjecture*. Geom. Topol. 11 (2007), 2117–2202. – cited on p. 1.

- [48] O. van Koert, *A Reeb flow on the three-sphere without a disk-like global surface of section*. Qual. Theory Dyn. Syst. 19, 36 (2020). – cited on p. 4.
- [49] E. Zehnder, *Homoclinic points near elliptic fixed points*, Comm. Pure Appl. Math. 26 (1973), 131 – 182. – cited on p. 27.

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