Some examples of global solutions associated with large initial data for the incompressible Navier-Stokes system

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Examples of large solutions to Navier-Stokes

Outline of the talk

The Cauchy problem

- Presentation of the equations
- Fundamental properties
- Weak solutions
- Strong solutions
- Towards the largest adapted space
- The Koch and Tataru space

Situations when large data can generate a global solution

- Previous results
- Two examples

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Presentation of the equations

- Viscous, incompressible, homogeneous fluid, in two or three space dimensions
- Velocity $u = (u^1, u^2, u^3)(t, x)$, pressure p(t, x)

(NS)
$$\partial_t u + u \cdot \nabla u - \Delta u = -\nabla p,$$

div $u = 0,$

with div
$$u = \sum_{j=1}^{3} \partial_{j} u^{j}$$
 and $u \cdot \nabla u = \sum_{j=1}^{3} u^{j} \partial_{j} u = \sum_{j=1}^{3} \partial_{j} (u^{j} u).$

Remark : The pressure can be eliminated by projection onto divergence-free vector fields :

$$\mathbb{P} = \mathsf{Id} -
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Cauchy data : $u_{|t=0} = u_0$.

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Fundamental properties

Conservation of the energy

$$\frac{1}{2}\frac{d}{dt}\|u(t)\|_{L^2}^2+\|\nabla u(t)\|_{L^2}^2=0$$

due to the structure of the nonlinear term : $(u \cdot \nabla u | u)_{L^2} = 0$.

Scale invariance If u is a solution of (*NS*) associated with the initial data u_0 on [0, *T*[, then for all $\lambda > 0$,

 $u_{\lambda}(t,x) \stackrel{\text{def}}{=} \lambda u(\lambda^2 t, \lambda x)$

is a solution associated with $u_{\lambda,0}(x) \stackrel{\text{def}}{=} \lambda u_0(\lambda x)$ on $[0, \lambda^{-2}T[$.

To solve (NS) one should try to use both informations.

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Theorem [Leray, 1934]

Let $u_0 \in L^2(\mathbb{R}^d)$ be a divergence free vector field. There is a solution u of (*NS*) satisfying for all $t \ge 0$

$$\|u(t)\|_{L^2}^2 + 2\int_0^t \|\nabla u(t')\|_{L^2}^2 dt' \le \|u_0\|_{L^2}^2.$$

In two space dimensions that solution is unique (and satisfies the energy equality).

Remarks :

- Proof by compactness.
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- Search for conditions on the initial data to guarantee uniqueness.

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Strong solutions

One does not use the structure of the equation, but rather its scale invariance, by a fixed point method.

Solving (*NS*) is equivalent to solving

 $u = \mathbb{S}(t)u_0 + \mathbb{B}(u, u)$

where $\mathbb{S}(t)$ is the heat semi-group on \mathbb{R}^d and \mathbb{B} the bilinear form

$$\mathbb{B}(u,u)(t)=-\int_0^t \mathbb{S}(t-t')\mathbb{P}\, ext{div}\,(u\otimes u)(t')\,dt'.$$

The problem consists in finding an adapted Banach space X, such that \mathbb{B} is continuous from $X \times X$ to X.

Remark : By the scale invariance, the norm on X must satisfy

 $\lambda \|f(\lambda^2 t, \lambda(x-a))\|_X \sim \|f\|_X.$

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An existence and uniqueness result

Theorem

Let X be an adapted space. If u_0 is such that $\|\mathbb{S}(t)u_0\|_X$ is small enough, then there is a unique solution to (NS) in X.

Proof: This is simply Picard's theorem : if X is a Banach space, if $L \in \mathcal{L}(X)$ and if $\mathbb{B} \in \mathcal{B}(X)$, with $\|L\|_{\mathcal{L}(X)} < 1$, then for all x_0 in X satisfying

$$\|x_0\|_X < \frac{(1-\|L\|_{\mathcal{L}(X)})^2}{4\|\mathbb{B}\|_{\mathcal{B}(X)}}$$

the equation

$$x = x_0 + Lx + \mathbb{B}(x, x)$$

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Some examples of adapted spaces with d = 3

- The space of functions u₀ such that t^{1/4}S(t)u₀ belongs to L[∞]([0, T]; H¹); this is guaranteed if u₀ is in H^{1/2} (Fujita-Kato, 1964).
- The space of functions u_0 such that $t^{\frac{1}{4}}\mathbb{S}(t)u_0$ belongs to $L^{\infty}([0, T]; L^6)$; this is guaranteed if u_0 is in L^3 (Kato 1983).
- The space of functions u_0 such that $t^{\frac{1}{2}(1-\frac{3}{p})}\mathbb{S}(t)u_0$ belongs to $L^{\infty}([0, T]; L^p)$, for p a real number in $]3, \infty[$; this corresponds to the Besov space $\dot{B}_{p,\infty}^{-1+\frac{3}{p}}$ (Cannone-Meyer-Planchon, 1993).

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Remarks

Those are small data or small time theorems.

They hold for the more general equation

 $\partial_t u - \Delta u + Q(u, u) = 0$

where $Q(v, w) \stackrel{\text{def}}{=} \sum_{1 \le j,k \le 3} Q_{j,k}(D)(v^j w^k)$ and $Q_{j,k}(D)$ are smooth homogeneous Fourier multipliers of order 1.

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Besov spaces and oscillations



Proof: If $t \leq \varepsilon^2$, one has $\|\mathbb{S}(t)f_{\varepsilon}\|_{L^p} \leq \|f\|_{L^p}$ hence $\|\|t^{\frac{\sigma}{2}}\mathbb{S}(t)f_{\varepsilon}\|_{L^p}\|_{L^1([0,\varepsilon^2],\frac{dt}{t})} \lesssim \varepsilon^{\sigma}\|f\|_{L^p}$

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If $t \ge \varepsilon^2$, since $-i\varepsilon \partial_3 e^{i\frac{x_3}{\varepsilon}} = e^{i\frac{x_3}{\varepsilon}}$, one has

$$\begin{split} \mathbb{S}(t)f_{\varepsilon}(x) &= -\varepsilon^{2}\int_{\mathbb{R}^{3}}\partial_{y_{3}}^{2}\left(G_{t}(x-y)f(y)\right)e^{i\frac{y_{3}}{\varepsilon}}dy\\ &= -\frac{\varepsilon^{2}}{t}\int_{\mathbb{R}^{3}}G_{t}(x-y)f(y)e^{i\frac{y_{3}}{\varepsilon}}dy\\ &+ 2\frac{\varepsilon^{2}}{t^{\frac{1}{2}}}\int_{\mathbb{R}^{3}}G_{t}(x-y)\partial_{3}f(y)e^{i\frac{y_{3}}{\varepsilon}}dy\\ &- \varepsilon^{2}\int_{\mathbb{R}^{3}}G_{t}(x-y)\partial_{3}^{2}f(y)e^{i\frac{y_{3}}{\varepsilon}}dy, \end{split}$$

where $G_t(x)$ is a function of the type $\frac{1}{(4\pi t)^{\frac{3}{2}}}e^{-\frac{|x|^2}{4t}}P(\frac{x}{\sqrt{t}})$, and the result follows from Young's inequality and integration in time.

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If $t \ge \varepsilon^2$, since $-i\varepsilon \partial_3 e^{i\frac{x_3}{\varepsilon}} = e^{i\frac{x_3}{\varepsilon}}$, one has

$$\begin{split} \mathbb{S}(t)f_{\varepsilon}(x) &= -\varepsilon^{2}\int_{\mathbb{R}^{3}}\partial_{y_{3}}^{2}\left(G_{t}(x-y)f(y)\right)e^{i\frac{y_{3}}{\varepsilon}}dy\\ &= -\frac{\varepsilon^{2}}{t}\int_{\mathbb{R}^{3}}G_{t}(x-y)f(y)e^{i\frac{y_{3}}{\varepsilon}}dy\\ &+ 2\frac{\varepsilon^{2}}{t^{\frac{1}{2}}}\int_{\mathbb{R}^{3}}G_{t}(x-y)\partial_{3}f(y)e^{i\frac{y_{3}}{\varepsilon}}dy\\ &- \varepsilon^{2}\int_{\mathbb{R}^{3}}G_{t}(x-y)\partial_{3}^{2}f(y)e^{i\frac{y_{3}}{\varepsilon}}dy, \end{split}$$

where $G_t(x)$ is a function of the type $\frac{1}{(4\pi t)^{\frac{3}{2}}}e^{-\frac{|x|^2}{4t}}P(\frac{x}{\sqrt{t}})$, and the result follows from Young's inequality and integration in time.

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Towards the largest adapted space

Proposition [Meyer]

Let B be a Banach space continuously embedded in $\mathcal{S}'(\mathbb{R}^3)$ and such that

 $\lambda \|f(\lambda(\cdot - a))\|_B \sim \|f\|_B.$

Then *B* is continuously embedded in $\dot{B}_{\infty,\infty}^{-1}$.

Proof : We notice that

 $|\langle f, e^{-|\cdot|^2} \rangle| \leq C ||f||_B$

By dilation and translation, we deduce that

 $\|f\|_{\dot{B}^{-1}_{\infty,\infty}} = \sup_{t>0} t^{\frac{1}{2}} \|\mathbb{S}(t)f\|_{L^{\infty}} \leq C \|f\|_{B}.$

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The Koch and Tataru space

Definition

We denote by X the space of functions f of $L^2_{loc}(\mathbb{R}^+ \times \mathbb{R}^3)$ such that

$$\|f\|_{X} \stackrel{\text{def}}{=} \sup_{t>0} \left(t^{\frac{1}{2}} \|f(t)\|_{L^{\infty}} + \sup_{\substack{x \in \mathbb{R}^{3} \\ R>0}} R^{-\frac{3}{2}} \left(\int_{P(x,R)} |f(t,y)|^{2} dy dt \right)^{\frac{1}{2}} \right) < \infty,$$

where $P(x, R) = [0, R^2] \times B(x, R)$.

Remark : The space of initial data u_0 such that $\mathbb{S}(t)u_0$ is in X is the space BMO^{-1} and we have $\dot{B}_{\infty,2}^{-1} \hookrightarrow BMO^{-1} \hookrightarrow \dot{B}_{\infty,\infty}^{-1}$.

Theorem [Koch, Tataru 2001] The operator \mathbb{B} is continuous from $X \times X$ to X.

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Link with the Montgomery-Smith toy Navier-Stokes model (Gallagher,Paicu 2008)

• Example 2 : slow variation in one direction (Chemin, Gallagher 2007)

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Theorem [Chemin, Gallagher 2006]

If u_0 is a smooth, divergence free vector field in \mathbb{R}^3 such that, for some (appropriate) space F

$$\|\mathbb{B}(\mathbb{S}(t)u_0,\mathbb{S}(t)u_0)\|_F \leq C \exp\left(-C\|u_0\|^4_{\dot{B}^{-1}_{\infty,2}}
ight),$$

then the system (NS) has a unique global solution.

• *F* is a space for forcing terms in the Koch-Tataru theorem.

• **Proof** : define $R(t) = u(t) - S(t)u_0$ then

 $R = \mathbb{B}(\mathbb{S}(t)u_0, \mathbb{S}(t)u_0) + 2\mathbb{B}(\mathbb{S}(t)u_0, R) + \mathbb{B}(R, R).$

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Example 1

Theorem [Chemin, Gallagher 2006] Let $\phi \in \mathcal{S}(\mathbb{R}^3)$ and $\alpha \in]0, 1[$. Set

$$\varphi_{\varepsilon}(x) = \frac{(-\log \varepsilon)^{\frac{1}{5}}}{\varepsilon^{1-\alpha}} \cos\left(\frac{x_3}{\varepsilon}\right) \phi\left(x_1, \frac{x_2}{\varepsilon^{\alpha}}, x_3\right).$$

There is C > 0 such that, if ε is small enough,

$$u_0^{\varepsilon}(x) \stackrel{\mathrm{def}}{=} (\partial_2 \varphi_{\varepsilon}(x), -\partial_1 \varphi_{\varepsilon}(x), 0)$$

satisfies

$$\mathcal{C}^{-1}(-\log \varepsilon)^{rac{1}{5}} \leq \|u_0^{\varepsilon}\|_{\dot{B}^{-1}_{\infty,\infty}} \leq \mathcal{C}(-\log \varepsilon)^{rac{1}{5}}, \quad \text{and}$$

 $\|\mathbb{B}(\mathbb{S}(t)u_0^{\varepsilon},\mathbb{S}(t)u_0^{\varepsilon})\|_F\leq Carepsilon^{rac{lpha}{3}}(-\logarepsilon)^{rac{2}{5}}.$

In particular for ε small enough, u_0^{ε} generates a unique, global, smooth solution to (*NS*).

The Montgomery-Smith toy model

The equation

$$(TNS_1) \quad \partial_t u - \partial_x^2 u + |\partial_x|(u^2) = 0$$

has the same small data theorems as Navier-Stokes, but some large data can generate finite-time blow up (Montgomery-Smith, 2001).

Theorem [Gallagher, Paicu 2008] Let d = 2 or 3. There is a bilinear operator Q, which is a d-dimensional matrix of Fourier multipliers, such as the equation $(TNS_{t}) \int \partial_{t} u - \Delta u + Q(u, u) = 0 \quad \text{in } \mathbb{R}^{+} \times \mathbb{R}^{d}$

$$u_{|t=0} = u_0$$

preserves the **divergence free condition** and the **scale invariance** of (NS), and such that there is a global, unique solution if the data is small enough in BMO^{-1} . Moreover the initial data of the previous theorem generates a finite-time blow up.

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Example 2 : slow variation in one direction

Theorem [Chemin, Gallagher 2007]

Let $v_0^h = (v_0^1, v_0^2)$ and w_0 be smooth divergence free vector fields on \mathbb{R}^3 . If ε is small enough, the initial data

 $u_0^{\varepsilon}(x) = \left(v_0^h + \varepsilon w_0^h, w_0^3\right)\left(x_h, \varepsilon x_3\right)$

generates a unique, global solution u^{ε} of (NS).

Remark : One needs to check that such an initial data exists (that is not small) : let (f,g) be in $S(\mathbb{R}^2)$ and $S(\mathbb{R})$ respectively. Let us consider the function h^{ε} defined by $h^{\varepsilon}(x_h, x_3) \stackrel{\text{def}}{=} f(x_h)g(\varepsilon x_3)$. We have, if ε is small enough,

 $\|h^{\varepsilon}\|_{\dot{B}^{-1}_{\infty,\infty}(\mathbb{R}^3)} \geq \frac{1}{4} \|f\|_{\dot{B}^{-1}_{\infty,\infty}(\mathbb{R}^2)} \|g\|_{L^{\infty}(\mathbb{R})}.$

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$$v_{app}^{\varepsilon}(t,x) \stackrel{\text{def}}{=} \left(\underline{v}^{h} + \varepsilon \underline{w}^{\varepsilon,h}, \underline{w}^{\varepsilon,3}\right)(t, x_{h}, \varepsilon x_{3})$$

and $\underline{v}^{h}(t, x_{h}, y_{3})$ solves the 2D Navier-Stokes system with data $v_{0}^{h}(x_{h}, y_{3})$ (y_{3} is a parameter) while $\underline{w}^{\varepsilon}$ solves a transport-diffusion equation

$$\partial_t \underline{w}^{\varepsilon} + \underline{v}^h \cdot \nabla_h \underline{w}^{\varepsilon} - \Delta_h \underline{w}^{\varepsilon} - \varepsilon^2 \partial_3^2 \underline{w}^{\varepsilon} = -(\nabla^h \underline{p}_1, \varepsilon^2 \partial_3 \underline{p}_1)$$

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Claim : R^{ε} is small, globally in time.

 $\partial_t R^{\varepsilon} + R^{\varepsilon} \cdot \nabla R^{\varepsilon} + v^{\varepsilon}_{app} \cdot \nabla R^{\varepsilon} + R^{\varepsilon} \cdot \nabla v^{\varepsilon}_{app} - \Delta R^{\varepsilon} = -\nabla q_{\varepsilon} + F^{\varepsilon},$ e $F^{\varepsilon}(t, x_h, x_3) \sim \varepsilon^{\frac{1}{3}}$ in $L^2(\mathbb{R}^+; \dot{H}^{-\frac{1}{2}}(\mathbb{R}^3)).$

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Indeed $F^{\varepsilon}(t, x_h, x_3)$ is made of the following terms :

$$\begin{split} \varepsilon \Big(\big(\varepsilon \underline{w}^{\varepsilon} \cdot \nabla \underline{w}^{\varepsilon,h}, \underline{w}^{\varepsilon} \cdot \nabla \underline{w}^{\varepsilon,3} \big) + \big(\underline{w}^{\varepsilon} \cdot \nabla \underline{v}^{h}, 0 \big) \Big) (t, x_{h}, \varepsilon x_{3}) \\ + \varepsilon \Big(\varepsilon \big(\partial_{3}^{2} \underline{v}^{h}, 0 \big) + \big(0, \partial_{3} \underline{p}_{0} \big) \Big) (t, x_{h}, \varepsilon x_{3}). \end{split}$$

Study of the nonlinear terms : For any smooth *a* and *b* and any $1 \le j \le 3$,

 $\begin{aligned} \|a\partial_j b\|_{L^2(\mathbb{R}^+;\dot{H}^{-\frac{1}{2}}(\mathbb{R}^3))} \lesssim \|a\partial_j b\|_{L^2(\mathbb{R}^+;L^{\frac{3}{2}}(\mathbb{R}^3))} \lesssim \|a\|_{L^\infty(\mathbb{R}^+;L^3(\mathbb{R}^3))} \|\partial_j b\|_{L^2(\mathbb{R}^+;L^3(\mathbb{R}^3))}. \end{aligned}$ This implies that

 $\|a\partial_j b(t,\mathsf{x}_h,\varepsilon\mathsf{x}_3)\|_{L^2(\mathbb{R}^+;\dot{H}^{-\frac{1}{2}}(\mathbb{R}^3))} \lesssim \varepsilon^{-\frac{2}{3}} \|a\|_{L^\infty(\mathbb{R}^+;\dot{H}^{\frac{1}{2}}(\mathbb{R}^3))} \|\partial_j b\|_{L^2(\mathbb{R}^+;\dot{H}^{\frac{1}{2}}(\mathbb{R}^3))}.$

Indeed $F^{\varepsilon}(t, x_h, x_3)$ is made of the following terms :

$$\begin{split} \varepsilon \Big(\big(\varepsilon \underline{w}^{\varepsilon} \cdot \nabla \underline{w}^{\varepsilon,h}, \underline{w}^{\varepsilon} \cdot \nabla \underline{w}^{\varepsilon,3} \big) + \big(\underline{w}^{\varepsilon} \cdot \nabla \underline{v}^{h}, 0 \big) \Big) (t, x_{h}, \varepsilon x_{3}) \\ + \varepsilon \Big(\varepsilon \big(\partial_{3}^{2} \underline{v}^{h}, 0 \big) + \big(0, \partial_{3} \underline{p}_{0} \big) \Big) (t, x_{h}, \varepsilon x_{3}). \end{split}$$

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3

The ill prepared case

Ideally one would like to consider initial data of the following type (with 0 $< \alpha \leq$ 1) :

$$u_0^{\varepsilon}(x) = \left(v_0^h + \varepsilon^{\alpha} w_0^h, \varepsilon^{\alpha-1} w_0^3\right) \left(x_h, \varepsilon x_3\right)$$

where $v_0^h = (v_0^1, v_0^2, 0)$ is a horizontal, smooth divergence free vector field on \mathbb{R}^3 and w_0 is a smooth divergence free vector field on \mathbb{R}^3 .

Work in progress (Chemin, Gallagher, Paicu 2008) : consider

$$u_0^{\varepsilon}(x) = \left(\varepsilon^{\frac{1}{2}}w_0^h, \varepsilon^{-\frac{1}{2}}w_0^3\right)(x_h, \varepsilon x_3)$$

in $\mathbb{T}^2 \times \mathbb{R}$ with $\int_{\mathbb{T}^2} w_0^h(x_h, x_3) \equiv 0$, with (w_0^h, w_0^3) small and smooth enough (analytic-type in x_3) then there is a global solution.

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